Asymptotically exponential hitting times and metastability:

a pathwise approach without reversibility

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Plan of the talk

- -1- Metastability and first hitting
- -2- Known results and tools
- -3- The non reversible case [FMNS]

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-4- Recurrence as a robust tool

-1- Metastability and first hitting

Physical systems near a phase transition (e.g. ferromagnet or saturated gas)

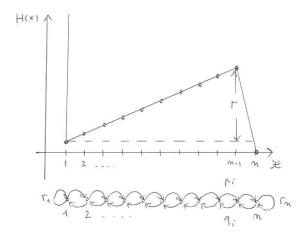
- trapped for an abnormal long time in a state —the metastable phase— different from the eventual equilibrium state consistent with the thermodynamical potentials;
- subsequently, undergoing a *sudden* transition at a *random time* from the metastable to the stable state.

Statistical mechanics model:

- space state \mathcal{X} , e.g. $\mathcal{X} = \{-1, +1\}^{\Lambda}$; interaction, e.g. Ising hamiltonian;
- evolution: Markov chain on \mathcal{X} , reversible w.r.t. Gibbs measure $\frac{e^{-\beta H}}{Z}$;
- decay of the metastable state: convergence to equilibrium of the chain, (equilibrium state), e.g. configuration of minimal energy in the limit of small temperature $\beta \rightarrow \infty$.

An example to introduce the problem

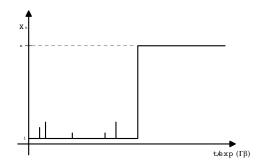
Random walk on $\mathcal{X} = \{1, ..., n\}$ reversible w.r.t. the following hamiltonians, i.e., with stationary measure $\pi(x) = \frac{e^{-\beta H(x)}}{z}$:



Let $\delta = H(x) - H(x - 1)$ for x = 2, ..., n - 1 and define , we have $\delta = 0$

$$p_x := P(x, x+1) = \frac{e^{-\delta\beta}}{1+e^{-\delta\beta}}, \quad x = 1, ..., n-1, \quad r_1 := P(1, 1) = 1-p_1,$$

$$q_x := P(x, x-1) = 1-p_x, \ x = 2, ..., n-1, \quad q_n = \frac{e^{-\beta [H(n-1)-H(n)]}}{1+e^{\delta \beta}}, \ r_n = 1-q_n$$



For large β (similarly large *n* [Barrera,Bertoncini,Fernandez]): after many unsuccessful attempts there is a fast transition to *n*.

Metastability is characterized by:

- Exit from a well (valley) of H with a motion against the drift: large deviation regime.
- Many visits to the metastable state (bottom of the well, point 1 in the ex.) before the transition to the stable one (n), large tunneling time with exponential distribution if properly rescaled.
- ► The existence of critical configurations separating the metastable state from the stable one, (n 1), first hitting to rare events.

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Our goal

Metastability / first hitting to rare events is usually studied in the literature for reversible Markov chain.

Our goal: prove asymptotic exponential behavior in a non reversible context, when $|\mathcal{X}| \to \infty$.

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Example: deck of *n* cards, $|\mathcal{X}| = n!$ Markov chain: top-in-at-random shuffling invariant maesure = uniform distribution first hitting to a particular configuration *G*.

- non reversible case
- entropic barrier

People:

...

The "first hitti	ng" community:
[K]	Keilson (1979) <i>(FM)</i>
[AB], [B]	Aldous, Brown (1982-92)
	Day (1983), Galves, Schmitt (1990),
[IMcD]	Iscoe, McDonald (1994),
	Asselah, Dai Pra (1997), Abadi, Galves (2001)
[FL]	Fill, Lyzinski (2012)

The "metastable" community:

[LP]	Lebowitz, Penrose (1966) <i>(FM)</i>
[FW]	Freidlin, Wentzell (1984) <i>(FM)</i>
[CGOV]	Cassandro, Galves, Olivieri, Vares (1984)
[it]	Martinelli, Olivieri, S. (1989)
[fr]	Catoni, Cerf (1995)
[B.et al]	Bovier, Eckhoff, Gayrard, Klein (2001)
[BL]	Bertrand, Landim (2011/12)
[BG]	Bianchi, Gaudillière (2012)

(FM) = founding member

-2- Known results and tools: first hitting community

The model: X_t ; $t \ge 0$ irreducible, finite-state, reversible Markov chain in continuous time, with transition rate matrix Q and stationary distribution π , so that $\pi Q = 0$ and

$$egin{aligned} \mathsf{DBC} : \pi_i \mathcal{Q}_{ij} &= \pi_j \mathcal{Q}_{ji} \ \mathcal{P} &= \mathbb{1} + \mathcal{Q} \end{aligned}$$

 $0 = \lambda_0 < \lambda_1 \le \lambda_2 \le ...$ real eigenvalues of the matrix -Q, $R = 1/\lambda_1$: relaxation time of the chain. If the set *A* is such that $R/\mathbb{E}_{\pi}\tau_A$ is small then is possible [AB] to obtain

estimate like :

$$|\mathbb{P}_{\pi}(\tau_{A}/\mathbb{E}_{\pi}\tau_{A} > t) - e^{-t}| \leq \frac{R/\mathbb{E}_{\pi}\tau_{A}}{1 + R/\mathbb{E}_{\pi}\tau_{A}} \qquad \forall t > 0$$
(1)

$$\mathbb{P}_{\pi}(\tau_{A} > t) \ge (1 - \frac{R}{\mathbb{E}_{\alpha}\tau_{A}}) \exp\{-\frac{t}{\mathbb{E}_{\alpha}\tau_{A}}\}$$
(2)

Moreover in the regime $R \ll t \ll \mathbb{E}_{\pi}\tau_A$ bounds on the density function are given [AB] in order to obtain a control on the distribution of τ_A also on scale smaller than $\mathbb{E}_{\pi}\tau_A$.

Results by metastability community

 $(\mathcal{X}^{(n)}, n \ge 1)$ sequence of finite state spaces, with $|\mathcal{X}^{(n)}| = n$ $(X_t^{(n)})_{t \in \mathbb{R}}$ sequence of continuous time, irreducible, reversible Markov chains on $\mathcal{X}^{(n)}$

 $Q^{(n)}$ transition rate matrix generating the chain $X_t^{(n)}$

 $(\pi^{(n)}, n \ge 1)$ invariant measures

asymptotics $n \to \infty$.

starting at $x \in \mathcal{X}^{(n)}$, and the hitting time to a set $G^{(n)} \subset \mathcal{X}^{(n)}$:

$$\tau_{G^{(n),x}}^{(n),x} = \inf \left\{ t \ge 0 : \ X_t^{(n),x} \in G^{(n)} \right\}$$
(3)

 $x_0^{(n)} \in \mathcal{X}^{(n)}$ = metastable state $G^{(n)} \subset \mathcal{X}^{(n)}$ = critical configurations (or stable state)

$$rac{ au_{G^{(n)}}^{x_0^{(n)}}}{\mathbb{E} au_{G^{(n)}}^{x_0^{(n)}}} \longrightarrow_{n o \infty}^{(d)} Y \sim \textit{Exp}(1)$$

Hypotheses for metastability

- i) pathwise approach: ([CGOV], [it], [fr]) recurrence to $x_0^{(n)}$ in a time $R_n \ll \mathbb{E}\tau_{G^{(n)}}^{(n)} x_0^{(n)}$ with large probability.
- ii) potential theoretical approach: ([B. et al], [BL])

$$Hp.A: \lim_{n \to \infty} n\rho_A(n) = 0$$
(4)
$$Hp.B: \lim_{n \to \infty} \rho_B(n) = 0$$
(5)

$$\rho_{A}(n) := \sup_{z \in \mathcal{X}^{(n)} \setminus \{x_{0}^{(n)}, G^{(n)}\}} \frac{\mathbb{P}\left(\tilde{\tau}_{G^{(n)}}^{(n), x_{0}^{(n)}} < \tilde{\tau}_{x_{0}^{(n)}}^{(n), x_{0}^{(n)}}\right)}{\mathbb{P}\left(\tilde{\tau}_{\{x_{0}^{(n)}, G^{(n)}\}}^{(n), z} < \tilde{\tau}_{z}^{(n), z}\right)}$$

$$\rho_{B}(n) := \sup_{z \in \mathcal{X}^{(n)} \setminus \{x_{0}^{(n)}, G^{(n)}\}} \frac{\frac{\mathbb{E}\tau_{\{x_{0}^{(n)}, G^{(n)}\}}^{(n), z}}{\mathbb{E}\tau_{G^{(n)}}^{(n), x_{0}^{(n)}}} \cdot \frac{\tilde{\tau}_{A}^{(n), x_{0}^{(n)}}}{\mathbb{E}\tau_{A}^{(n), x}} = \min\left\{t > 0 : X_{t}^{(n), x} \in A\right\}$$

Some tools (fh)

- collapsed chain technique: hitting to a single state $A \equiv j$ ([K], [AB])
- For reversible chains for each *j*, by using spectral representation and Laplace transform, τ_j^{π} is a geometric convolution of suitable i.i.d.r.v. W_i :

$$\tau_j^{\pi} = \sum_{i=1}^{N} W_i \tag{6}$$

with *N* a geometric random variable of parameter π_j , approximately exponential in the sup norm [B] when $\pi(j)$ is small.

- [FL] generalize (6) to non reversible chains under additional hypotheses.
- τ_i^{π} is completely monotone [K]:

$$\mathbb{P}_{\pi}(\tau_j > t) = \sum_{l=1}^{m} p_l \exp\{-\gamma_l t\}$$
(7)

with $p_i \ge 0$ and $0 < \gamma_1 < \gamma_2 < ... < \gamma_m$ the distinct eigenvalues of $-Q_j$. Complete monotonicity is a powerful tool and exponential behavior follows from it [AB].

- interlacing between eigenvalues of $-Q_i$ and -Q

$$\mathbf{0} = \lambda_{\mathbf{0}} < \gamma_{\mathbf{1}} \leq \lambda_{\mathbf{1}} \leq \gamma_{\mathbf{2}} \leq \lambda_{\mathbf{2}} \leq \dots$$

Canceling out common pairs of eigenvalues from the two spectra and renumbering them we obtain

$$\mathbf{0} = \lambda_{\mathbf{0}} < \gamma_{\mathbf{1}} < \lambda_{\mathbf{1}} < \gamma_{\mathbf{2}} < \lambda_{\mathbf{2}} < \ldots < \gamma_{m} < \lambda_{m}$$

again by Laplace transform [B]

$$au_j^{\pi} \sim \sum_{i=1}^m Y_i$$

with

$$\mathbb{P}(Y_i > t) = \left(1 - \frac{\gamma_i}{\lambda_i}\right) e^{-\gamma_i t}, \qquad t > 0, \ i = 1, ..., m$$

Moreover

$$\left(1-\frac{\gamma_1}{\lambda_1}\right)e^{-\gamma_1 t} \leq \mathbb{P}_{\pi}(\tau_j > t) \leq (1-\pi(j))e^{-\gamma_1 t}$$

- [IMcD] exponential behavior in the non reversible case under additional implicit hypotheses (related to auxiliary processes involved in the proof) by studying the smallest real eigenvalue of a suitable Dirichlet problem.

Some tools (m)

- Different strategies in different regimes.
 Freidlin Wentzell techniques, cycle decomposition, cycle paths...
 In FW theory reversibility is not required, cycles and cycle path are easily defined in the reversible case.
- ii) Spectral characteristic of the generator, tools from potential theory, variational principles,...

$$c(i,j)=\frac{1}{r(i,j)}=\pi_i P_{ij}$$

$$r(i,j) = r(j,i) \iff \mathsf{DBC}$$

Extension to non reversible chains by Eckhoff (not published) (??) [BL] non reversible case under additional implicit hypotheses, which are not easy to verify in the non reversible case.

-3-The non reversible case [FMNS1]

 $(\mathcal{X}^{(n)}, n \ge 1)$ sequence of finite state spaces, with $|\mathcal{X}^{(n)}| = n$ $(\mathcal{X}_t^{(n)})_{t \in \mathbb{R}}$ continuous time, irreducible Markov chains on $\mathcal{X}^{(n)}$ $x_0^{(n)}$ = metastable state, $G^{(n)}$ = critical configurations (or stable state) in the following sense:

Hp. $G(T_n)$:

there exist sequences $r_n \rightarrow 0$ and $R_n \ll T_n$ such that

$$\sup_{x\in\mathcal{X}}\mathbb{P}\Big(\tau^{(n),x}_{\{x_0^{(n)},G^{(n)}\}}>R_n\Big) \leq r_n.$$

asymptotic results: the following are equivalent

$$\frac{\tau_{G^{(n)}}^{x_0^{(n)}}}{T_n} \longrightarrow_{n \to \infty}^{(d)} Y \sim Exp(1)$$

$$\exists \ell \geq 1, \xi < 1 \ : \ P\Big(\tau_{G^{(n)}}^{(n), x_0^{(n)}} > \ell \ T_n\Big) \leq \xi \text{ uniformly in } n$$

quantitative results:

$$|\mathbb{P}_x(\tau_G/\mathbb{E}_{x_0}\tau_G>t)-e^{-t}|\leq f(\frac{R}{E_{x_0}\tau_G},r)$$

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Comparison of hypotheses

[FMNS]

$$T_n^{LT} =$$

mean local time spent in $x_0^{(n)}$ before reaching $G^{(n)}$ starting from $x_0^{(n)}$

$$T_n^{G^{\xi}} := \inf\left\{t : \mathbb{P}\left(\tau_{G^{(n)}}^{(n), x_0^{(n)}} > t\right) \le \xi\right\}$$

$$T_n^{\mathsf{E}} = \mathbb{E}\left(\tau_{G^{(n)}}^{(n), x_0^{(n)}}\right)$$

The following implications hold:

 $Hp.A \implies Hp.G(T_n^{LT}) \implies Hp.G(T_n^{E}) \iff Hp.(T_n^{Q^{\xi}}) \iff Hp.B$

for any $\xi < 1$. Furthermore, the missing implications are false.

-4-Recurrence as a robust tool

a) Factorization property:

If S > R with $\sup_{x\in\mathcal{X}}\mathbb{P}\Big(\tau^x_{\{x_0,G\}}>R\Big) \leq r.$ then for any $t, s > \frac{R}{s}$ $\mathbb{P}ig(au_G^{x_0} > (t+s)Sig) \stackrel{\geq}{=} ig[\mathbb{P}ig(au_G^{x_0} > tS + Rig) - rig] \mathbb{P}ig(au_G^{x_0} > sSig) \ \leq ig[\mathbb{P}ig(au_G^{x_0} > tS - Rig) + rig] \mathbb{P}ig(au_G^{x_0} > sSig) \;.$ $\tau(tS) = \inf\{T > tS; X_T \in \{x_0, G\}\}$ $\mathbb{P}(\tau(tS) - tS > R) < r$

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b) Control on $\mathbb{P}(\tau_G^{x_0} \leq tS \pm R)$

Let
$$S > R$$
 with $\sup_{x \in \mathcal{X}} \mathbb{P}\Big(au_{\{x_0,G\}} > R\Big) \leq r$.

then

$$\begin{array}{ll} \mathbb{P}(\tau_G^{x_0} \leq S) \leq & \mathbb{P}(\tau_G^{x_0} \leq S + R) & \leq & \mathbb{P}(\tau_G^{x_0} \leq S) \left[1 + a\right] & (8) \\ \mathbb{P}(\tau_G^{x_0} \leq S) \geq & \mathbb{P}(\tau_G^{x_0} \leq S - R) & \geq & \mathbb{P}(\tau_G^{x_0} \leq S) \left[1 - a\right] & (9) \end{array}$$

with

$$a = \frac{\mathbb{P}(\tau_G^{x_0} \le 2R)}{\mathbb{P}(\tau_G^{x_0} \le S)} + \frac{r}{\mathbb{P}(\tau_G^{x_0} \le S)}$$
 (10)

This result is useful when $\mathbb{P}(\tau_G^{x_0} \leq 2R)$ and *r* are small w.r.t. $\mathbb{P}(\tau_G^{x_0} \leq S)$. In this case we can conclude that *a* is small and so we get a multiplicative error estimate.

c) Exponentian behavior

By iterating a) and b) we get the exponential law on a suitable time scale in the interval $(R, \mathbb{E}\tau_G^{x_0})$. Let $T := \mathbb{E}\tau_G^{x_0}$ and $\varepsilon := \frac{R}{T}$

$$b = \mathbb{P}(\tau_G^{x_0} \le 2R) + 2r. \tag{11}$$

If $\varepsilon + r < \frac{1}{4}$ and $S < T(1 - 4(\varepsilon + r))$. Then $\mathbb{P}(\tau_G^{x_0} > S) > b$ and for each positive integer k

$$\mathbb{P}(au_G^{ extsf{x}_0} > k oldsymbol{S}) \stackrel{\leq}{=} \left[\mathbb{P}(au_G^{ extsf{x}_0} > oldsymbol{S}) + oldsymbol{b}
ight]^k \ \geq \ \left[\mathbb{P}(au_G^{ extsf{x}_0} > oldsymbol{S}) - oldsymbol{b}
ight]^k$$

d) Generic starting point

By recurrence to $\{x_0, G\}$ we have for all $x \in B(x_0)$

$$\tau_G^{\mathbf{X}} \sim \tau_G^{\mathbf{X}_0}$$

Improvement of recurrence

These result are relevant only if *r* is small. It is possible to decrease exponentially *r* by increasing linearly $\varepsilon := \frac{R}{T}$ with $T = \mathbb{E}\tau_G^{\chi_0}$. Let *R* be such that

$$\sup_{x\in\mathcal{X}}\mathbb{P}\Big(\tau^x_{\{x_0,G\}}>R\Big) \leq r.$$

and suppose $\varepsilon = \frac{R}{T}$ small.

We can chose another return time $R^+ \in (R, T)$ with $\Gamma := \frac{R^+}{R} < \frac{1}{\varepsilon}$. Define $\varepsilon^+ := \frac{R^+}{T} = \varepsilon \Gamma < 1$. The recurrence property in time R^+ is immediately estimate by the following

$$\sup_{x \in \mathcal{X}} \mathbb{P}(\tau_{\{x_0, G\}}^x > R^+) \le r^{\frac{R^+}{R}} = r^{\Gamma} =: r^+$$
(12)

This meas that with this new recurrence time R^+ we have $\varepsilon^+ = \varepsilon \Gamma$ and $r^+ = r^{\Gamma}$. From this we get a control on exponential behavior on small time scales

Time scale for exponential behavior

With this improvement of recurrence, if

$$\lim_{n\to\infty}\frac{(r_n)^{\frac{R_n^+}{R_n}}}{p_n} = 0.$$

with

$$p_n := \mathbb{P}\Big(au_{G^{(n)}}^{(n),x_0^{(n)}} < 3R_n^+\Big) \to 0$$

and if S_n is such that $R_n^+ < S_n \le T_n$, then $\tau_{G^{(n)}}^{(n), x_0^{(n)}}$ has asymptotic exponential behaviour at scale S_n i.e., for every integer k

$$\lim_{n \to \infty} \frac{\mathbb{P}(\tau_{G^{(n)}}^{(n), x_0^{(n)}} \in (kS_n, (k+1)S_n])}{\mathbb{P}(\tau_{G^{(n)}}^{(n), x_0^{(n)}} > S_n)^k \ \mathbb{P}(\tau_{G^{(n)}}^{(n), x_0^{(n)}} \le S_n)} = 1$$

Recurrence on a set instead of on a single state x_0 [FMNS2]

Work in progress.

Back to the example: deck of n cards, top-in-at-random. First hitting to a particular configuration G.

- ► mixing time of order T_{mix} = n log n ⇒ recurrence with large probability to a suitable set B of configurations such that:
 - B is "large" : π(B) > 1 o_n(1);
 - τ_G^x is controlled uniformely in *B*:

$$\sup_{x\in B}\mathbb{P}(\tau_G^x < T_{mix}) \leq f_n \to 0.$$

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• recurrence in $B \implies$ asymptotic exponential behavior.