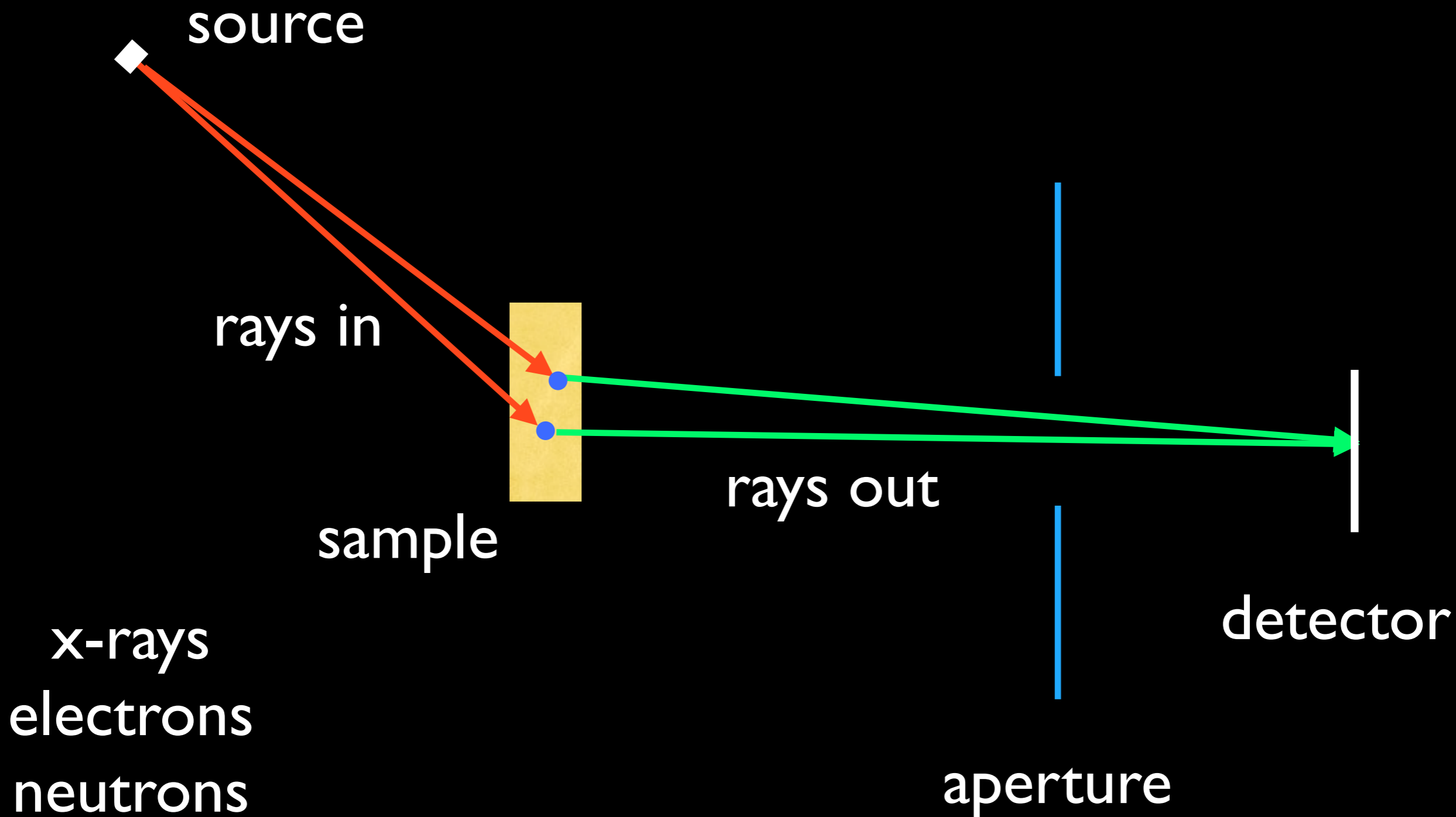


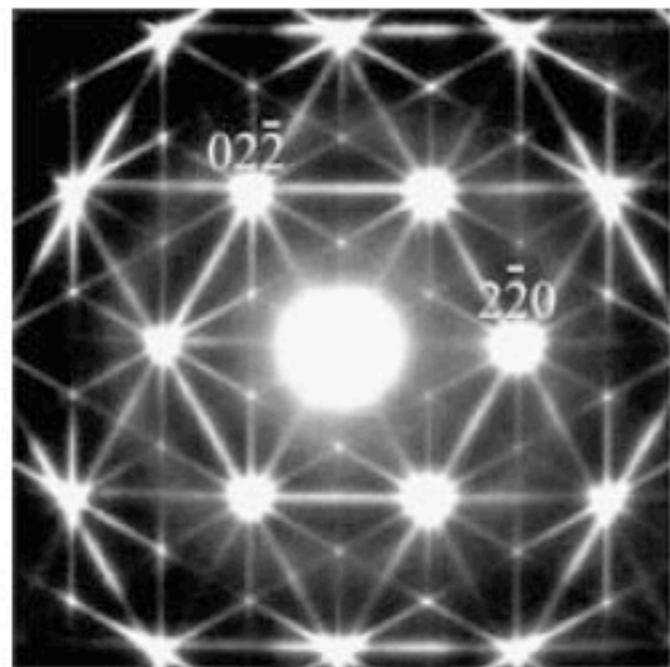
**Stationary spatial stochastic  
processes  
and the inverse problem  
in pure point diffraction**

**Robert V. Moody, University of Victoria**

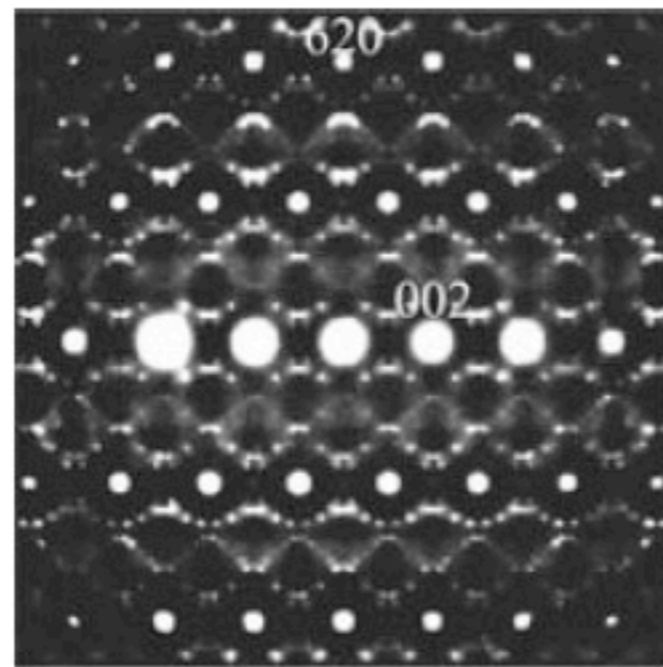
**BIRS, September 2011**



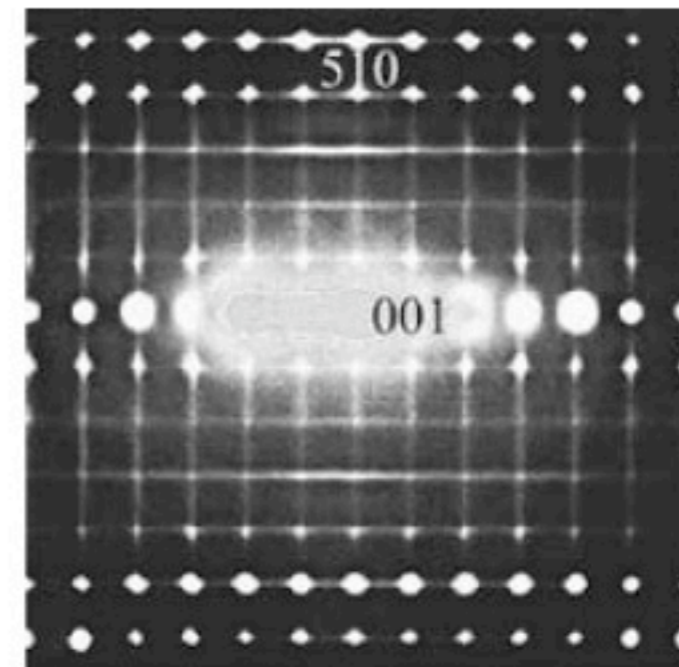
# Microscope/Diffractometer



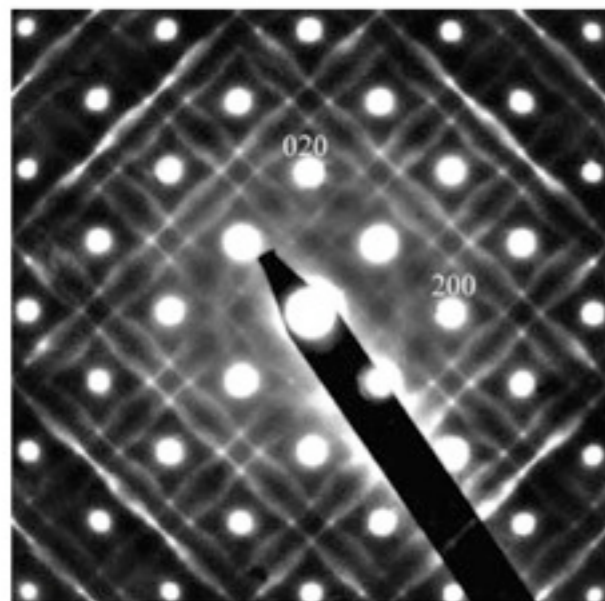
(a)



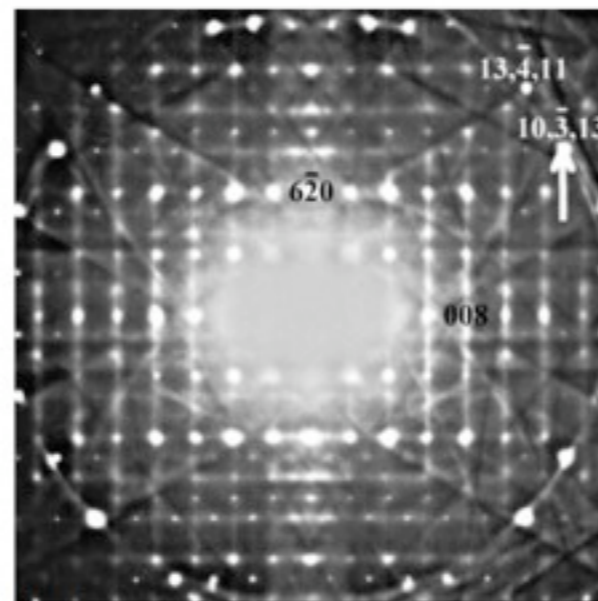
(b)



(c)



(a)



(b)

**FIGURE 4** An (a) [001] zone axis EDP of  $P4/nmm$  ThAsSe (see Withers *et al.*, 2006 for details) and (b) a  $\langle 130 \rangle$  zone axis electron diffraction pattern of  $Im\bar{3}$  CCTO (see Liu *et al.*, 2005 for details). Note that the  $\mathbf{G} \pm \sim 0.14 \langle 110 \rangle^* \pm \epsilon \langle 1, -1, 0 \rangle^*$  diffuse streaking in (a) occurs only around the  $h + k$  odd,  $n$ -glide forbidden parent reflections, whereas the diffuse streaking perpendicular to  $\langle 001 \rangle$  in (b) only runs through the  $[hkl]^*$ ,  $l$  even parent reflections.

from R. Withers:  
 Disorder: structured  
 diffuse scattering and  
 local crystal  
 chemistry

# Diffraction

---

self-interference of scattered radiation  
from a scattering density or distribution of  
matter  $\rho$

$\rho$

- could be a set of point masses
- the electric field of an assemblage of atoms
- more generally a measure in  $\mathbb{R}^d$

Mathematically the diffraction is

$$\omega = |\hat{\rho}|^2$$

treated as a measure in  $\mathbb{R}^d$

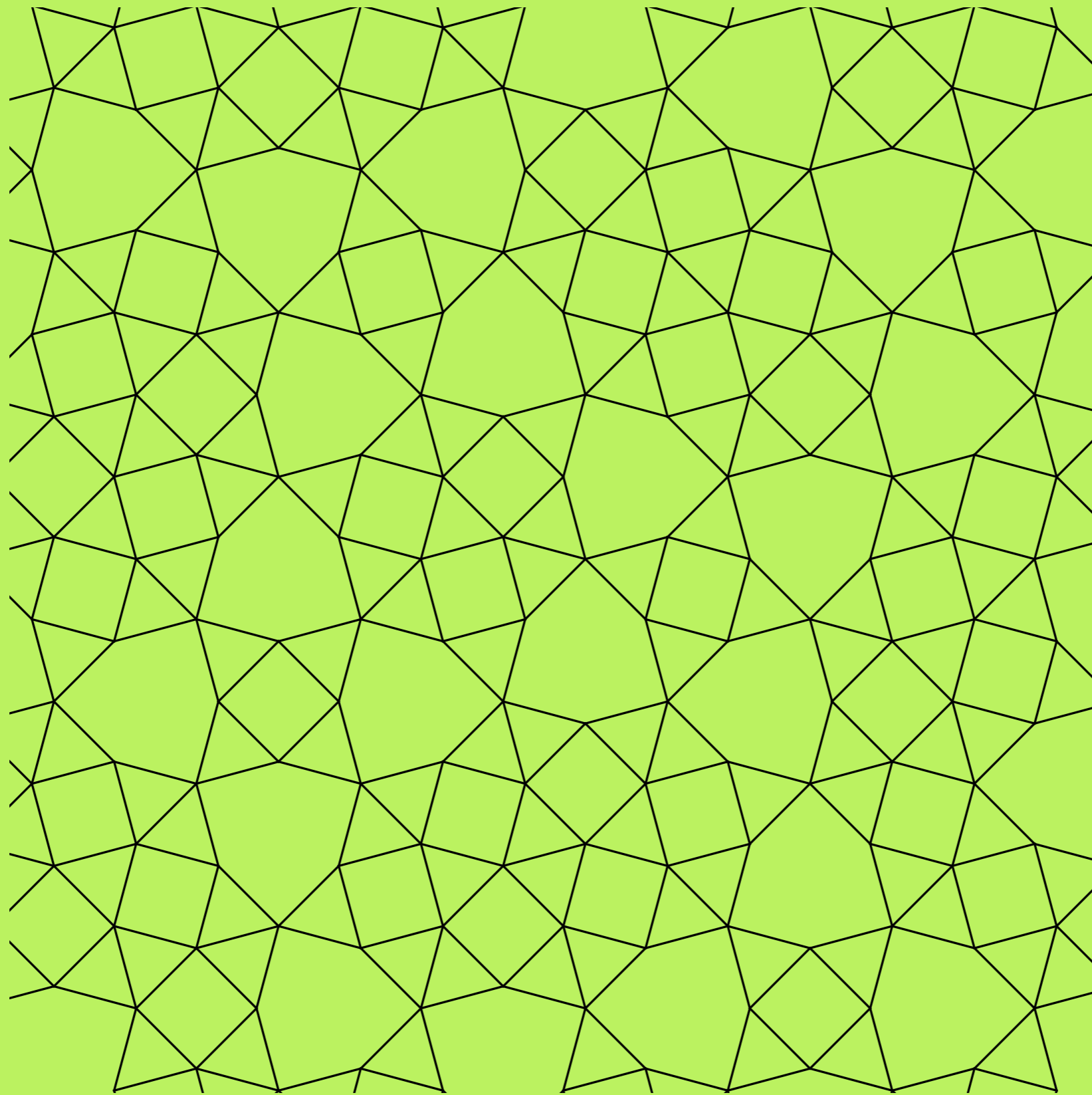
The diffraction is a measure in reciprocal space -- the Fourier dual of the physical space

In the case of  $\mathbb{R}^d$  this is another copy of  $\mathbb{R}^d$ .

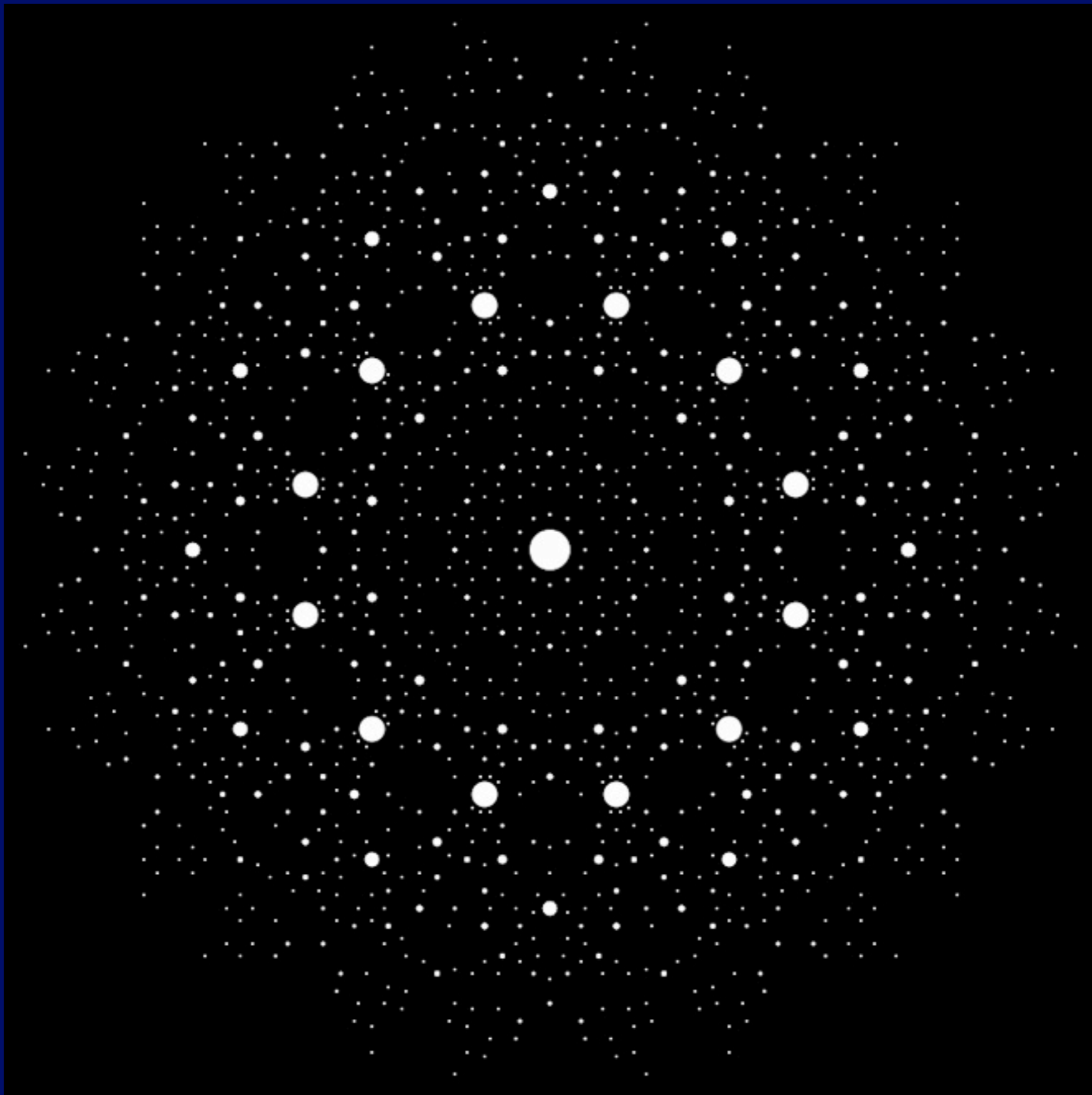
Later  $\mathbb{R}^d$  will be replaced by any locally compact Abelian group  $G$ ,  
Fourier space becomes  $\widehat{G}$ .

## Three basic properties of the diffraction measure

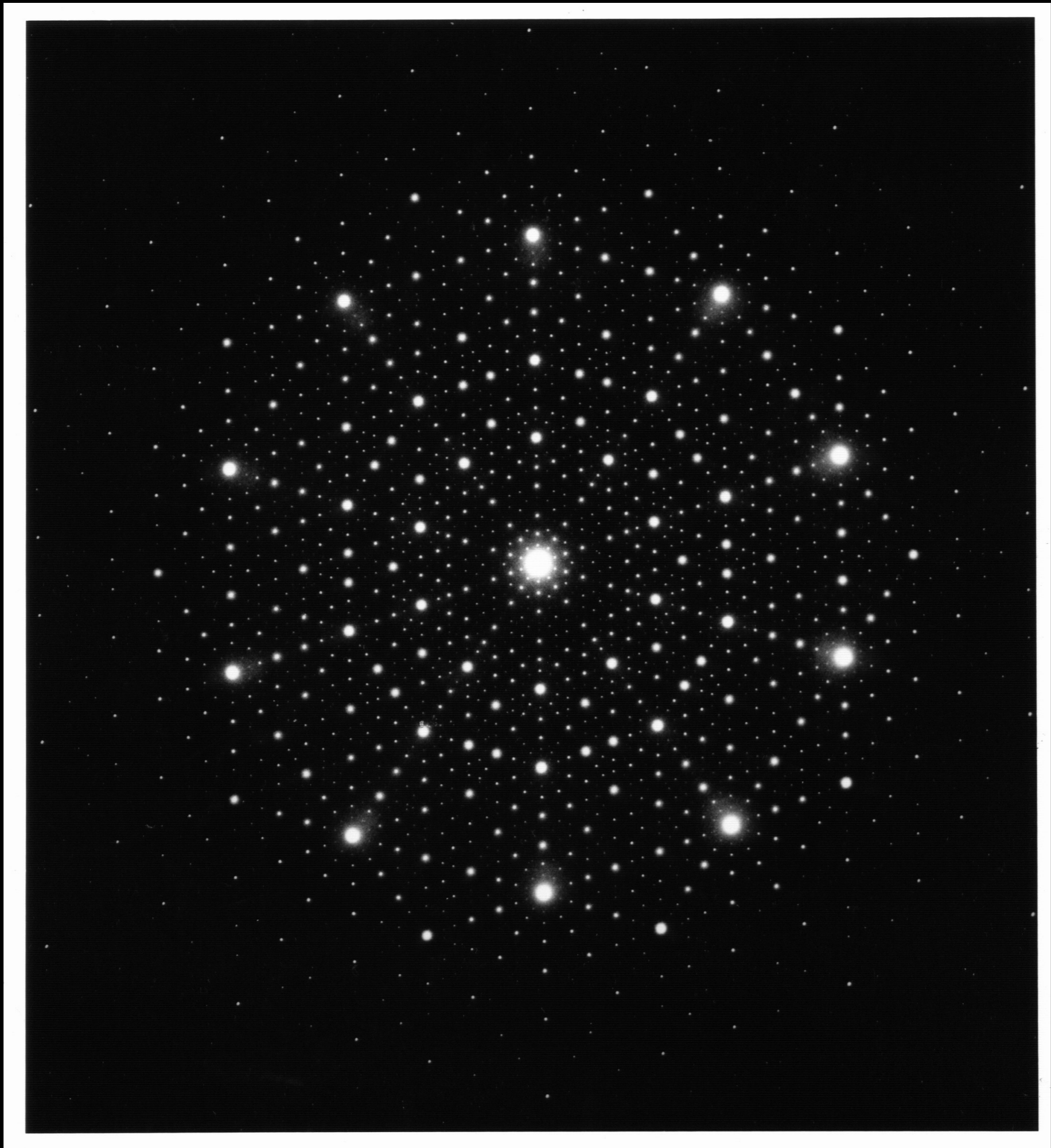
- It is positive
- It is centrally symmetric
- It is a translation bounded measure



## The Shield Tiling -- Franz Gaehler



# Diffraction of the Shield Tiling





the diffraction is a positive and centrally symmetric measure in Fourier space

The central problem of diffraction theory:

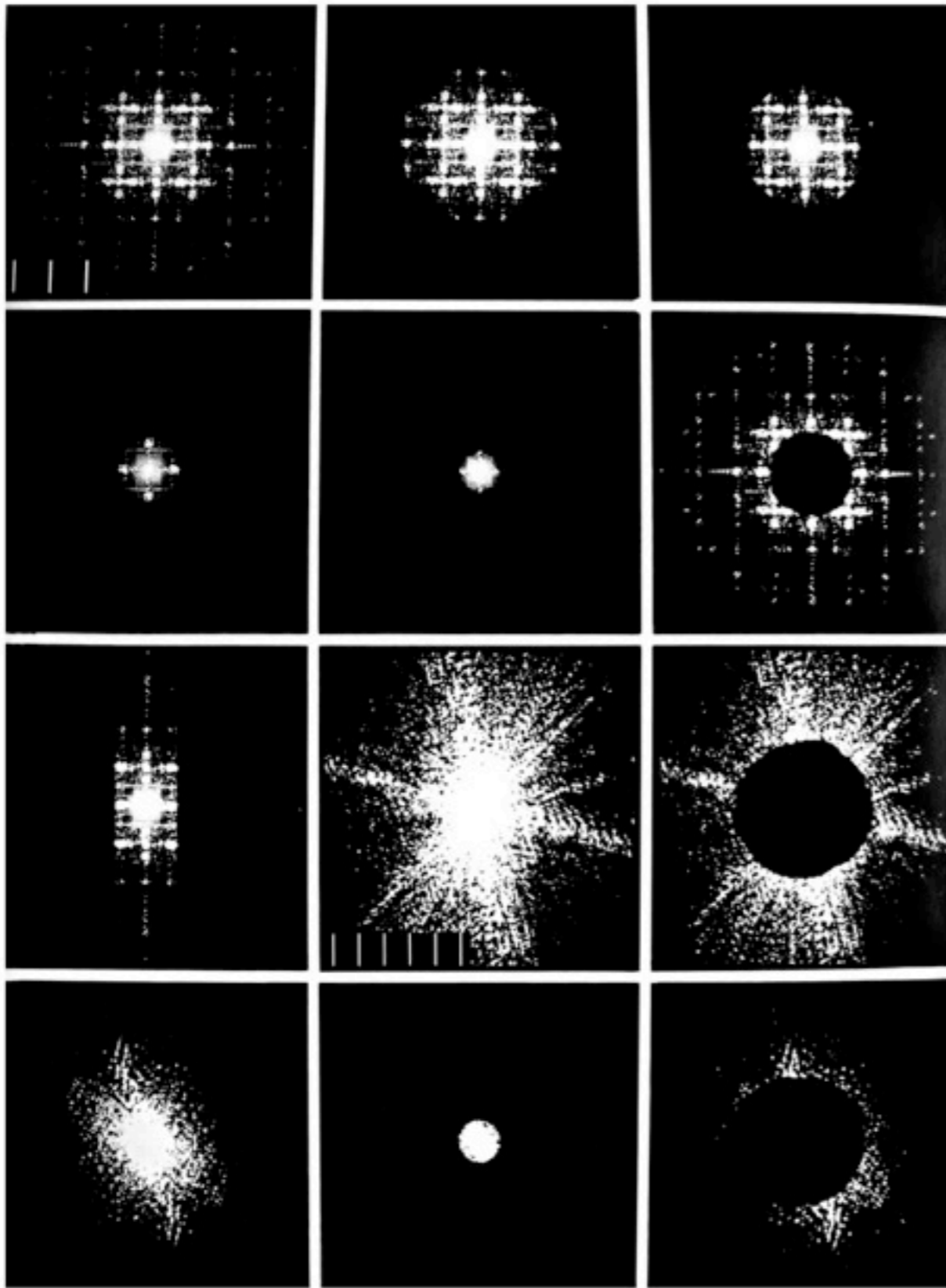
given the diffraction, determine the structure that created it

Why it is hard: We want  $\rho$

We are given  $\omega = |\hat{\rho}|^2$  This is the Phase Problem

$\omega = \hat{\rho} \overline{\hat{\rho}}$  the Patterson  
"function"

$= \widehat{\rho * \tilde{\rho}} = \hat{\gamma}$   a.k.a.  
the autocorrelation



# Brief History

---

Knipping/Laue experiment 1912


Bragg and Bragg (William Henry and William Lawrence)  
1913 first explanations

Arthur (Lindo) Patterson function 1934

Linus Pauling 1939 bixbyite – homometry

D. Shechtman, quasicrystals, 1982

different structures  
with the same  
diffraction



# Three problems

1. Given a centrally symmetric positive translation bounded measure, find a structure which has it as its diffraction.
2. Classify all the structures with the same diffraction.
3. What do we mean by a “structure”?

In the pure point case we can give reasonable answers to these three questions

A key point is to introduce the notion of stationary stochastic processes

This is the main subject of this talk

joint work with Daniel Lenz, Jena, Germany

# stationary spatial stochastic process

---

$$N : C_c(\mathbb{R}^d) \longrightarrow L^2(X, \mu)$$
$$F \longmapsto N(F)$$

$$C_c(\mathbb{R}^d)$$

the space of continuous functions of compact support

$$(X, \mu) = (X, \mathbb{R}^d, \mu)$$

a measure dynamical system with an  $\mathbb{R}^d$ -action on it

$\mu$  is stationary (translation invariant)

$N$  is a linear  $\mathbb{R}^d$  - map

$N$  is real: real valued functions go to real valued functions

# motivation from stochastic point processes

---

a random variable whose outcome for each event is  
a discrete point set in  $\mathbb{R}^d$

$X$  is the set of all possible outcomes.

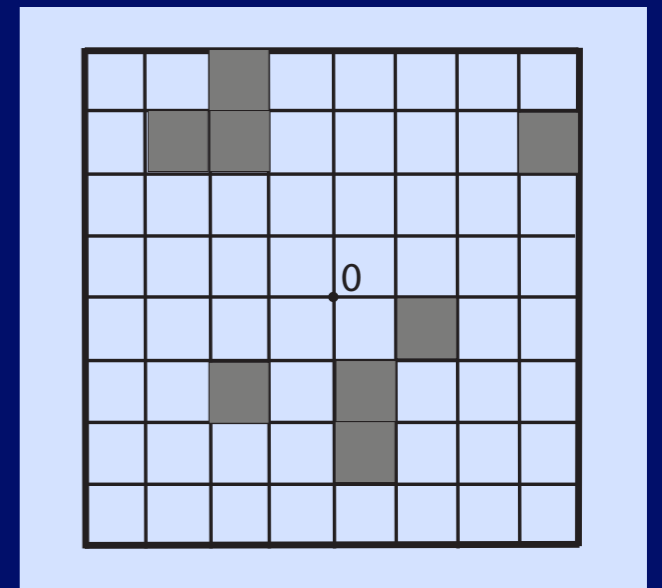
each  $\xi \in X$  is a discrete point set in  $\mathbb{R}^d$ .

$X$  is assumed to be  $\mathbb{R}^d$ -invariant

## measurable sets

given a pattern,  $\Pi$

$B(\Pi) =$  all  $\Lambda \in X$  matching the pattern  $\Pi$



these generate our sigma-algebra of measurable sets

$X$  is a stationary probability space

probability, or frequency measure  $\mu$

$\mu(B(\Pi)) =$  frequency of occurrence of the pattern  $\Pi$

$\left\{ \begin{array}{l} \mu \text{ is translation invariant (stationary)} \\ \mu(X) = 1 \end{array} \right.$

$\implies (X, \mathbb{R}^d, \mu)$

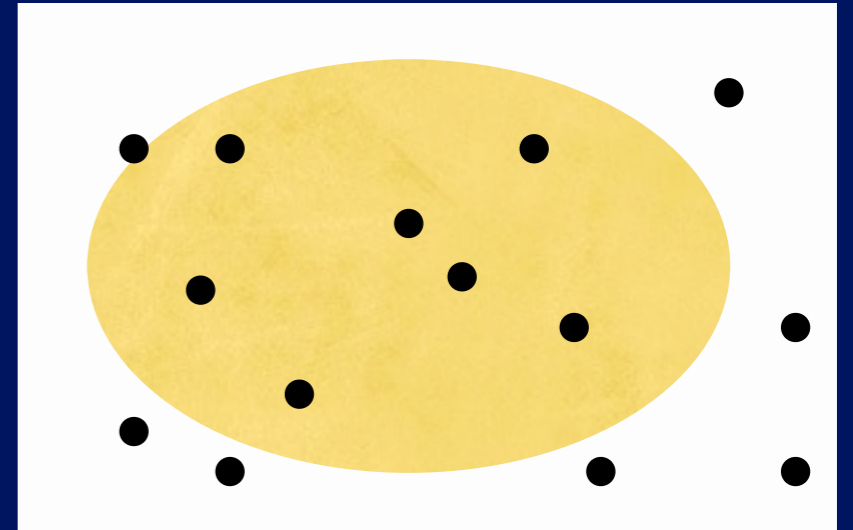
probability space  
measure theoretic dynamical system

**Assume the system is ergodic**



## defining N

For each measurable set  $A \subset \mathbb{R}^d$   
and for each  $\xi \in X$



$$\begin{aligned} N_A(\xi) &= \text{the number of points of } \xi \text{ in } A \\ &= \sum_{x \in \xi} 1_A(x) \end{aligned}$$

$$N_A : X \longrightarrow \mathbb{R} \quad \text{counting functions}$$

In effect, N provides a link between the physical space of the objects which are labelled by the elements of  $X$

extend to functions

---

$F$  function on  $\mathbb{R}^d$   $\longrightarrow$   $N(F)$  function on  $X$

$$N(F)(\xi) = \sum_{x \in \xi} F(x)$$

$$N : C_c(\mathbb{R}^d) \longrightarrow L^2(X, \mu)$$

stationary spatial stochastic processes (sssp)

We shall always assume that  $N$  has a second moment (to come soon!)

$$(X, \mathbb{R}^d, \mu, N)$$

the hull:  
compact  
space

physical  
space

probability  
measure

counting  
function

spatial stationary stochastic process

$$N : C_c(\mathbb{R}^d) \longrightarrow L^2(X, \mu)$$

$$F \mapsto N(F)$$

# how to think about N

---

$$N : C_c(\mathbb{R}^d) \longrightarrow L^2(X, \mu)$$

$$F \mapsto N(F)$$

Think of N as providing a way of providing test functions for the various elements of X

$$\xi : F \mapsto N(F)(\xi) \text{ for a.e. } \xi \in X.$$

Comment

locally compact Abelian  
**G** group

**H** Hilbert space

~~$$N : C_c(\mathbb{R}^d) \longrightarrow L^2(X, \mu)$$~~

~~$$F \mapsto N(F)$$~~

spatial  
versus  
non-spatial

two types: sssp and ssp

## Second moments

---

$N$  has a second moment if  
for all real-valued  $F, G \in C_c(\mathbb{R}^d)$

$$\mu^{(2)}(F \otimes G) = \langle N(F) | N(G) \rangle$$

is a measure on  $C_c(\mathbb{R}^d \times \mathbb{R}^d)$



there exists a positive definite measure  $\gamma$  on  $\mathbb{R}^d$  so that

$$\int_{\mathbb{R}^d} F * \tilde{F} d\gamma = \langle N(F) | N(F) \rangle$$

$\gamma$  is the autocorrelation

Its Fourier transform  $\omega = \hat{\gamma}$  is the diffraction

$\omega$  is positive, centrally symmetric, and translation bounded.

The question: given

$\omega$  is positive, centrally symmetric,  
and translation bounded.

Does there exist a stationary spatial stochastic  
process  
with this as its diffraction?

Yes, in the pure point case

pure point diffraction means  $\omega$  is a pure point measure.

Theorem:

In the ergodic case this is equivalent to  $L^2(X, \mu)$   
is a pure point dynamical system

# Some results

---

$N$  has a second moment iff  
for all real-valued  $F, G \in C_c(\mathbb{R}^d)$   
with support in the compact set  $K$   
there is a constant  $C_K$  so that  
$$\langle N(F) | N(G) \rangle \leq C_K \max\{ \|F\|_2 \|G\|_2, \|F\|_\infty \|G\|_1 \}$$

---

Any stationary spatial stochastic process  $N$  (sssp)  
with diffraction  $\omega$  splits into  $N = N_p + N_c$  where  
 $N_p$  is a pure point sssp with diffraction  $\omega_p$   
 $N_c$  is a continuous ssp with diffraction  $\omega_c$ .

---

A sssp  $N$  has pure point diffraction if and only if  
for all  $F \in C_c(\mathbb{R}^d)$ ,  $t \mapsto T_t N(F)$   
is strongly almost periodic with respect to the  
 $L^2$  norm on  $L^2(X, \mu)$ .

How does one construct a sssp from a given pure point positive centrally symmetric translation bounded measure?

---

Start with  $\omega = \sum_{k \in S} \omega(k) \delta_k$

$S$  is the set of Bragg peaks

There are two ways to approach this

Both need the concept of a phase form

$a : S \longrightarrow U(1)$  (the unit circle in  $\mathbb{C}$ )

$a(0) = 1$

$a(-k) = \overline{a(k)}$  for all  $k \in S$ .



1. Since  $L(X, \mu)$  should be pure point, Halmos-von Neumann says that measure theoretically  $X$  is a compact Abelian group.

$$E := \langle S \rangle \subset \mathbb{R}^d$$

is the subgroup generated by the Bragg peaks.

Give it the discrete topology and put

$$X = \widehat{E}.$$

This is a compact Abelian group  
and with usual Haar measure serves as  $(X, \mu)$ .

**the phase form enters in the definition of N**

## 2. Use Gel'fand theory

form a commutative Banach algebra  
and make  $X$  its set of maximal ideals

for  $k \in S$ , put  $c(k) = a(k) \omega(k)^{1/2}$

$\mathcal{P}$  is defined as the free commutative algebra over  $\mathbb{C}$   
with generators  $f_k, k \in S$

relations  $f_{k_1} \cdots f_{k_m} = c(k_1) \cdots c(k_m)$

whenever  $k_1 + \cdots + k_m = 0$ .

action of  $\mathbb{R}^d$

$$t \cdot f_k = (k, t) f_k$$

inner product

$$\langle f_{k_1} \cdots f_{k_m} | f_{l_1} \cdots f_{l_n} \rangle = \begin{cases} c(k_1) \cdots c(k_m) c(-l_1) \cdots c(-l_n) \\ 0 \end{cases}$$

according as  $k_1 + \cdots + k_m = l_1 + \cdots + l_n$  or not.

Use this to define the operator norm  $\nu$  on  $\mathbb{P}$ :  
 $\nu(x)$  is the norm of multiplication by  $x$ .

$\mathcal{A}$  is the completion of  $\mathcal{P}$  under  $\nu$ .

$\mathcal{A}$  is a commutative Banach algebra

$X$  is its maximal ideal spectrum.

skip the definition of the probability measure  $\mu$

$$\begin{aligned} N : C_c(\mathbb{R}^d) &\longrightarrow L^2(x, \mu) \\ F &\longmapsto \sum_{k \in S} \hat{F}(k) f_k \end{aligned}$$

this is the desired sssp with diffraction  $\omega$ .

$$(X, \mathbb{R}^d, \mu, N)$$

the hull:  
compact  
space

physical  
space

probability  
measure

counting  
function

spatial stationary stochastic process

$$N : C_c(\mathbb{R}^d) \longrightarrow L^2(X, \mu)$$

$$F \mapsto N(F)$$

The remaining problem is to classify phase forms  
and decide when they lead to equivalent sssp's

But what does the density really look like?

# Example

periodic structure on  $\mathbb{R}$  with diffraction  $\sum_{k \in \mathbb{Z}} \delta_k$

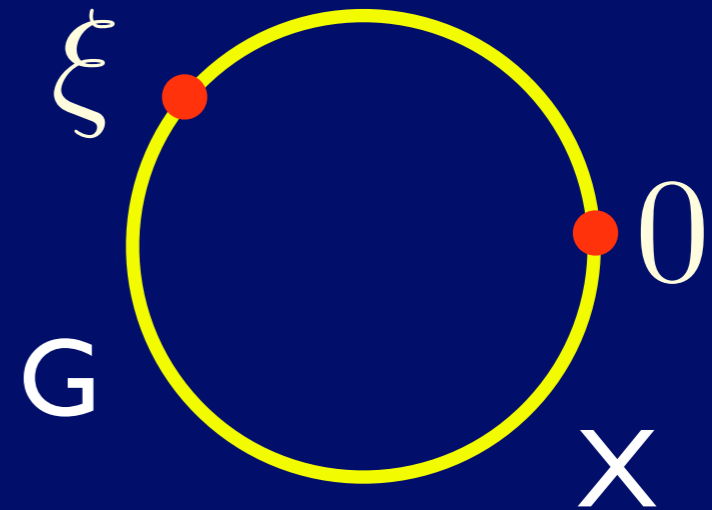
$$G = U(1) \quad \widehat{G} = \mathbb{Z} \quad X \simeq U(1)$$

$$\omega = \sum_{k \in \mathbb{Z}} \delta_k$$

$$\begin{aligned} N_a(F) &= \sum_{k \in \mathbb{Z}} \widehat{F}(k) a(k) \chi_k \\ &= \sum_{k \in \mathbb{Z}} \widehat{F}(k) a(k) e^{2\pi i k \cdot (\cdot)} \end{aligned}$$

If  $a \equiv 1$  then

$$\begin{aligned} N_1(F)(\xi) &= \sum_{k \in \mathbb{Z}} \widehat{F}(k) \chi_k(\xi) \\ &= \sum_{k \in \mathbb{Z}} \widehat{F}(k) e^{2\pi i k \cdot (\xi)} = F(\xi) = \delta_\xi(F) \end{aligned}$$



$$N_1 \Leftrightarrow \rho = \delta_0$$

What does it mean that  $N_1(F)(\xi) = \delta_\xi(F)$ ?

$N_1$  is supposed to offer some sort of ‘density’ at each point  $\xi_0 \in X$ .

$F \mapsto N_1(F)(\xi)$  shows how test functions interact with the density

In our case,  $X$  is just one orbit – that of  $1 \in U(1)$ .

From translation invariance of  $N$

we expect

density at  $\xi$  = translation by  $\xi$  of the density at 0.

$$N_1 \leftrightarrow \delta_0$$

## Example continued

---

Fix a finite set  $K = -K \subset \mathbb{Z} \setminus \{0\}$ .

Define phase form  $a$  by

$$a(k) = \begin{cases} 1, & \text{if } k \notin K \\ \text{arbitrary in } U(1), & \text{otherwise.} \end{cases}$$

with  $a(-k) = \overline{a(k)}$ .

Then  $\rho = \delta_0 + \sum_{k \in K} (a(k) - 1)e^{-2\pi i k \cdot (\cdot)}$ .

Remarkably this has the same diffraction,

$$\omega = \sum_{k \in \mathbb{Z}} \delta_k$$

# Moments

the  $m$  th moment

$$\mu^{(m)}(F_1 \otimes \cdots \otimes F_m) = \int N(F_1) \cdots N(F_m) d\mu$$

the second moment determines diffraction

knowing all moments determines  $\mu$

Example (Grunbaum and Moore)

$$G = \mathbb{Z}/6\mathbb{Z}, \quad \hat{G} = \frac{1}{6}\mathbb{Z}/\mathbb{Z}$$

$$\rho_1 = 11\delta_0 + 25\delta_1 + 42\delta_2 + 45\delta_3 + 31\delta_4 + 14\delta_5$$

$$\rho_2 = 10\delta_0 + 17\delta_1 + 35\delta_2 + 46\delta_3 + 39\delta_4 + 21\delta_5$$

these two densities have the same their first  
5 moments equal!!



# Summary

---

We can explicitly define sssp for any pure point measure that looks like a diffraction measure

We can classify all such sssp's

The notion of an sssp is very general and offers a good context for studying diffraction

The most outstanding problem is that we do not know if sssp's are always due to some sort of measure on the physical space.

We have no constructive theory for continuous or mixed diffraction types

Charles Radin  
Stephen Dworkin  
Bert Hof  
Martin Schlottmann  
Jean Bellissard  
Jean-Baptiste Gouéré  
Michael Baake  
Daniel Lenz  
Nicolae Strungaru

