Linearization and local existence of solutions for the volume preserving mean curvature flow with line tension

Lars Müller

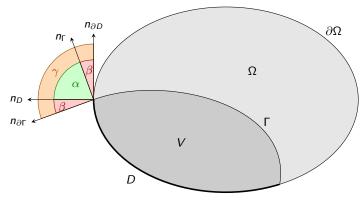
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General situation and notation



Basic technical assumption: $0 < \alpha(p) < \pi$ for all $p \in \partial \Gamma$.

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Energy functional

We consider the energy functional

$$E(\Gamma) := \int_{\Gamma} 1 d\mathcal{H}^2 - a \int_{D} 1 d\mathcal{H}^2 + b \int_{\partial \Gamma} 1 d\mathcal{H}^1 + \lambda \left(\int_{V} 1 dx - V_0 \right)$$

for $a, b, V_0 \in \mathbb{R}$ with $b \ge 0$.



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for $a, b, V_0 \in \mathbb{R}$ with $b \ge 0$. Varying the hypersuface by

$$\psi:\mathbb{R} imes\mathbb{R}^3\longrightarrow\mathbb{R}^3:(t,p)\longmapsto\psi(t,p):=p+t\zeta(p)$$

where $\zeta : \mathbb{R}^3 \longrightarrow \mathbb{R}^3$ is a tangential vectorfield to D and obtaining a family of evolving hypersurfaces by $\Gamma(t) := \psi(t, \Gamma)$, gives the first variation

$$\frac{d}{dt}E(\Gamma(t))\Big|_{t=0} = \int_{\Gamma} (\lambda - H_{\Gamma})(n_{\Gamma} \cdot \zeta) d\mathcal{H}^{2} + \int_{\partial \Gamma} (n_{\partial \Gamma} - an_{\partial D} - b\vec{\varkappa}) \cdot \zeta d\mathcal{H}^{1}$$

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This shows that a minimizer of that energy necessarily satisfies

$$H_{\Gamma} = \lambda = const.$$
 on Γ

$$0 = a + b \varkappa_{\partial D} + \langle n_{\Gamma}, n_{D} \rangle \qquad \text{on } \partial \Gamma$$

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One possible flow that tends towards a minimizer of the energy E is given by

$$V_{\Gamma} = H_{\Gamma} - \overline{H_{\Gamma}}$$
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where $\overline{H_{\Gamma}}$ is the mean of the mean curvature given by

$$\overline{H_{\Gamma}}:=rac{1}{\int\limits_{\Gamma}1d\mathcal{H}^{2}}\int_{\Gamma}H_{\Gamma}d\mathcal{H}^{2}$$



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Additionally there are several reasonable choices of boundary conditions. We will impose the boundary condition

$$v_{\partial D} = a + b \varkappa_{\partial D} + \langle n_{\Gamma}, n_{D}
angle$$
 on $\partial \Gamma$

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We introduce the coordinate system Ψ over a fixed reference hypersurface Γ^* as $\Psi : \Gamma^* \times (-\varepsilon_0, \varepsilon_0) \longrightarrow \Omega : (q, w) \longmapsto \Psi(q, w) := q + wn_{\Gamma^*}(q) + t(q, w)T(q)$

where $T : \Gamma^* \longrightarrow \mathbb{R}^3$ is an arbitrary tangential vectorfield, that coincides with $n_{\partial\Gamma^*}$ on $\partial\Gamma^*$ and vanishes outside a small neighborhood of $\partial\Gamma^*$ and

$$t: \Gamma^* imes (-arepsilon_0, arepsilon_0) \longrightarrow \mathbb{R}: (q, w) \longmapsto t(q, w)$$

is some smooth function such that $\Psi(q, 0) = q$ for all $q \in \Gamma^*$ and $\Psi(q, w) \in \partial \Omega$ for all $q \in \partial \Gamma^*$ and all $w \in (-\varepsilon_0, \varepsilon_0)$.

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With the help of the curvilinear coordinate system Ψ one can write the evolving hypersurface as a familiy of graphs over the fixed hypersurface Γ^* . To this purpose define a distance function

$$arrho:\mathbb{R}_+ imes \mathsf{\Gamma}^* \longrightarrow (-arepsilon_0,arepsilon_0):(t,q)\longmapsto arrho(t,q)$$

and set $\Gamma_{\varrho}(t) := \operatorname{Im}(\Psi(\bullet, \varrho(t, \bullet)))$ and observe that by construction one has $\Gamma_{\varrho\equiv 0}(t) = \Gamma^*$ for all $t \in \mathbb{R}_+$.

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With this notation the flow from above transforms into

$$\begin{split} V_{\Gamma_{\varrho}(t)}(\Psi(q,\varrho(t,q))) &= H_{\Gamma_{\varrho}(t)}(\Psi(q,\varrho(t,q))) - \overline{H}(\varrho(t)) & \text{in } \Gamma^* \\ v_{\partial D_{\varrho}(t)}(\Psi(q,\varrho(t,q))) &= a + b\varkappa_{\partial D_{\varrho}(t)}(\Psi(q,\varrho(t,q))) \\ &+ \left\langle n_{\Gamma_{\varrho}(t)}(\Psi(q,\varrho(t,q))), n_{D_{\varrho}(t)}(\Psi(q,\varrho(t,q))) \right\rangle & \text{on } \partial \Gamma^* \end{split}$$

which is a non-linear second order PDE in ϱ with second order boundary conditions.

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Problem

This PDE is also non-local due to $\overline{H}(\varrho(t))!$

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Linearization

The linearization of the PDE reads as

$$\begin{split} \partial_{t}\varrho(t) &= \Delta_{\Gamma^{*}}\varrho(t) + |\sigma^{*}|^{2}\varrho(t) + (\nabla_{\Gamma^{*}}H_{\Gamma^{*}} \cdot P\left(\partial_{w}\Psi(0)\right))\varrho(t) \\ &- \int_{\Gamma^{*}} \left(\Delta_{\Gamma^{*}} + |\sigma^{*}|^{2} - H_{\Gamma^{*}}^{2} + \overline{H}(\mathbb{O})H_{\Gamma^{*}})\varrho(t)d\mathcal{H}^{2} \\ &+ \frac{1}{\int_{\Gamma^{*}} 1d\mathcal{H}^{2}} \int_{\partial\Gamma^{*}} \left(H_{\Gamma^{*}} - \overline{H}(\mathbb{O})\right)\cot(\alpha)\varrho(t)d\mathcal{H}^{1} \quad \text{ on } \Gamma^{*} \\ \partial_{t}\varrho(t) &= -\sin(\alpha)^{2}(n_{\partial\Gamma^{*}} \cdot \nabla_{\Gamma^{*}}\varrho(t)) - \sin(\alpha)H_{D^{*}}(n_{\partial D^{*}}, n_{\partial D^{*}})\varrho(t) \\ &+ \sin(\alpha)\cos(\alpha)H_{\Gamma^{*}}(n_{\partial\Gamma^{*}}, n_{\partial\Gamma^{*}})\varrho(t) + b\sin(\alpha)\varrho_{\sigma\sigma}(t) \\ &+ b\sin(\alpha)\varkappa_{D^{*}}H_{D^{*}}(n_{\partial D^{*}}, n_{\partial D^{*}})\varrho(t) - b\sin(\alpha)\varkappa_{\partial D^{*}}\langle \vec{\tau}^{*}, (n_{\partial D^{*}})_{\sigma}\rangle \varrho(t) \\ &- b\sin(\alpha)\langle n_{\partial D^{*}}, (n_{D^{*}})_{\sigma}\rangle^{2}\varrho(t) \quad \text{ on } \partial\Gamma^{*} \\ \varrho(0) &= \varrho_{0} \quad \text{ on } \Gamma^{*} \end{split}$$

Highest order terms Non-local terms

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