# INVARIANT MEASURES OF GENERALIZED STOCHASTIC POROUS MEDIUM EQUATIONS 

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In this paper, we extend the results of the recent works [1], [2] on existence of infinitesimally invariant measures for the stochastic porous medium equation. The corresponding partial differential equation (see [3]) is

$$
\begin{equation*}
\partial x(u, t) / \partial t=\Delta \Psi(x)(u, t)+\Phi(x)(u, t), \tag{1}
\end{equation*}
$$

where $\Psi$ and $\Phi$ are certain functions, e.g., polynomials. The associated stochastic partial differential equation is heuristically written as

$$
\begin{equation*}
d x_{t}=\sqrt{2} d W_{t}+\left[\Delta \Psi\left(x_{t}\right)+\Phi\left(x_{t}\right)\right] d t . \tag{2}
\end{equation*}
$$

However, the rigorous interpretation is not obvious in case of nonlinear functions $\Psi$ and $\Phi$. One of the possible approaches to this problem is to consider the associated infinite dimensional elliptic operator $L$ on a suitable domain, find an infinitesimally invariant measure $\mu$ for $L$, and construct a Markovian semigroup on $L^{2}(\mu)$ having $\mu$ as in invariant measure such that the generator of the semigroup extends $L$, and finally construct a strong Markov process with continuous paths that solves the martingale problem corresponding to (1). In the case $\Psi(s)=$ $s^{m}+\alpha s, \alpha>0$, and $\Phi=0$, where $m$ is an odd number, this programme has been fulfilled in [2] and existence of an infinitesimally invariant measure for $\alpha=0$ and $m=3$ has been proved in [1]. Here we consider more general $\Psi$ and nonzero $\Phi$. In addition to greater generality of our assumptions, a novelty of this paper is that it provides constructive finite dimensional approximations of the invariant measure. The existence result is an application of a result of our earlier work [4].

Let $D \subset \mathbb{R}^{d}$ be a bounded open domain with a smooth boundary and let $\left\{e_{n}\right\}$ be the orthonormal basis in $L^{2}(D)$ formed by the eigenfunctions of the Laplacian $\Delta$ with Dirichlet boundary conditions. Thus, $\Delta e_{i}=\lambda_{i} e_{i}$ and we assume that $\lambda_{1} \leq \lambda_{2} \leq \cdots$. The inner product and norm in $L^{2}(D)$ are denoted by $(x, y)_{2}$ and $\|x\|_{2}$. Let $H_{0}^{2,1}(D)$ be the classical Sobolev space obtained as the completion of $C_{0}^{\infty}(D)$ with respect to the Sobolev norm $\|x\|_{2,1}=\|x\|_{2}+\||\nabla x|\|_{2}$.

Let $r>1$ and let $\zeta_{r}(s):=|s|^{r} \operatorname{sgn} s, s \in \mathbb{R}^{1}$. If $r$ is an odd number, then $\zeta_{r}(s)=s^{r}$. Let $\Psi$ be a $C^{1}$-function with $\Psi(0)=0$ such that for some positive numbers $\kappa_{0}, C_{0}$, and $\kappa_{1}$ one has

$$
\kappa_{0}|s|^{r-1} \leq \Psi^{\prime}(s) \leq C_{0}+\kappa_{1}|s|^{r-1} \quad \text { for all } s \in \mathbb{R}^{1},
$$

and let $\Phi$ be a continuous function satisfying the following condition: $|\Phi(s)| \leq C+\delta|s|^{r}$, where $0<\delta<4 \kappa_{0} \lambda_{1}(r+1)^{-2}$ and $C$ is a constant. For example, it suffices that $|\Phi(s)| \leq \kappa_{2}+\kappa_{3}|s|^{q}$, where $q \in(0, r), \kappa_{2}, \kappa_{3} \in(0,+\infty)$.

[^0]We study the existence of infinitesimally invariant measures for the infinite dimensional elliptic operator $L$ informally given by

$$
L f:=\Delta_{Q} f+\langle b, \nabla f\rangle, \quad b(x)=\Delta \Psi(x)+\Phi(x)
$$

on smooth cylindrical functions defined on $X:=L^{2}(D)$ or on the negative Sobolev space $H:=H^{2,-1}(D)$. A rigorous interpretation is this. Let $\mathcal{F} \mathcal{C}_{0}^{\infty}$ be the linear span of the class of all functions $f$ on $X$ of the form $f(x)=f_{0}\left(x_{1}, \ldots, x_{n}\right), f_{0} \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right), x_{i}=\left(x, e_{i}\right)_{2}$. Let

$$
b_{i}(x):=\int_{D}\left[\Psi(x(u)) \Delta e_{i}(u)+\Phi(x(u)) e_{i}(u)\right] d u, \quad x \in L^{r}(D) .
$$

Let $q_{i}>0$ be such that $S:=\sum_{i=1}^{\infty} q_{i}<\infty$. The operator

$$
L f:=\sum_{i=1}^{\infty}\left[q_{i} \partial_{e_{i}}^{2} f+b_{i} \partial_{e_{i}} f\right], \quad f \in \mathcal{F} \mathcal{F}_{0}^{\infty},
$$

where $\partial_{e_{i}}$ stands for the partial derivative along $e_{i}$, is well defined (for every $f \in \mathcal{F} \mathcal{C}_{0}^{\infty}, L f$ is just a finite sum). The second order part of $L$ can be regarded as trace $\left(Q D^{2} f\right)$, where $Q$ is the operator on $X$ defined by $Q e_{i}=q_{i} e_{i}$. The operator $Q$ is the covariance of the Wiener process $W_{t}$ in (2).

We shall say that a Borel probability measure $\mu$ on $X$ is infinitesimally invariant for $L$ if $\mu\left(L^{r}(D)\right)=1, b_{i} \in L^{1}(\mu)$ for all $i$ and

$$
\begin{equation*}
\int_{X} L f(x) \mu(d x)=0 \quad \forall f \in \mathcal{F} \mathcal{C}_{0}^{\infty} \tag{3}
\end{equation*}
$$

We write this symbolically as $L^{*} \mu=0$.
Let $E_{n}$ denote the linear space spanned by $e_{1}, \ldots, e_{n}$ and let

$$
L_{n} f:=\sum_{i=1}^{n}\left[q_{i} \partial_{e_{i}}^{2} f+b_{i} \partial_{e_{i}} f\right] .
$$

The orthogonal projection in $L^{2}(D)$ to $E_{n}$ is denoted by $P_{n}$.
Lemma 1. Let $\psi$ be a $C^{1}$-function on the real line such that $\psi(0)=0$.
(i) Suppose that $\left|\psi^{\prime}(s)\right| \geq C|s|^{r-1}$, where $C>0$ and $r \geq 1$. Let $f \in L^{2}(D)$ and $\psi \circ f \in$ $H_{0}^{2,1}(D)$. Then $|f|^{r} \operatorname{sgn} f \in H_{0}^{2,1}(D)$.
(ii) Suppose that $\left|\psi^{\prime}(s)\right| \leq C^{\prime}|s|^{r-1}$, where $C^{\prime}>0$ and $r \geq 1$. Let $f \in L^{2}(D)$ and $|f|^{r} \operatorname{sgn} f \in$ $H_{0}^{2,1}(D)$. Then $\psi \circ f \in H_{0}^{2,1}(D)$.

Proof. (i) The inverse function to $\psi$ will be denoted by $\eta$. Then $\zeta_{r} \circ x=\zeta_{r} \circ \eta \circ \psi \circ x$. The function $\eta$ is continuous, strictly increasing and differentiable outside the origin. Now it suffices to observe that the function $h:=\zeta_{r} \circ \eta$ is Lipschitzian and $h(0)=0$. Indeed,

$$
\left|h^{\prime}(s)\right|=\left|\zeta_{r}^{\prime}(\eta(s)) \eta^{\prime}(s)\right|=\left|\zeta_{r}^{\prime}(\eta(s)) / \psi^{\prime}(\eta(s))\right| \leq r / C
$$

Assertion (ii) is proved analogously.
We observe that the assumption $\psi(0)=0$ is only needed to ensure the zero boundary condition; in the case of $H^{2,1}(D)$, the same reasoning applies without that assumption. In place of the continuous differentiability of $\psi$ one can require that it is Lipschitzian (then the estimate on $\left|\psi^{\prime}\right|$ should hold a.e.).

Note that the inclusion $\zeta_{r} \circ x \in H_{0}^{2,1}(D)$ implies the inclusion $|x|^{r} \in H_{0}^{2,1}(D)$, but obviously is not equivalent to the latter.

Theorem 1. (i) Under the above assumptions, there exists a Borel probability measure $\mu$ on $X$ that is infinitesimally invariant for $L$ and is concentrated on the set of functions $x$ such that $\zeta_{(r+1) / 2} \circ x \in H_{0}^{2,1}(D)$ and

$$
\begin{equation*}
\int_{X} \int_{D}\left|\nabla\left(\zeta_{(r+1) / 2} \circ x\right)(u)\right|^{2} d u \mu(d x)<\infty \tag{4}
\end{equation*}
$$

(ii) If, in addition, $\sum_{i=1}^{\infty} q_{i} \sup _{u \in D} e_{i}(u)^{2}=: M<\infty$, then there exists a Borel probability $\mu$ that is infinitesimally invariant for $L$ and concentrated on the set of functions $x$ such that $\Psi \circ x \in$ $H_{0}^{2,1}(D)$ and consequently $\zeta_{r} \circ x \in H_{0}^{2,1}(D)$ and one has

$$
\begin{equation*}
\int_{X} \int_{D}|\nabla(\Psi \circ x)(u)|^{2} d u \mu(d x)<\infty \tag{5}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{X} \int_{D}\left|\nabla\left(\zeta_{r} \circ x\right)(u)\right|^{2} d u \mu(d x)<\infty . \tag{6}
\end{equation*}
$$

Finally, (4) holds. So (6) remains valid for $\zeta_{s}$ in place of $\zeta_{r}$ with any s between $(r+1) / 2$ and $r$.
Proof. (i) We verify that the hypotheses of Theorem 5.1 in [4] are satisfied. Those hypotheses are:
(a) existence of functions $V: X \rightarrow[0,+\infty]$ and $\Theta: X \rightarrow[0,+\infty]$ such that the sets $\{\Theta \leq c\}$ are compact and the restrictions of $V$ to the subspaces $E_{n}$ are compact $C^{2}$-functions,
(b) the continuity of the functions $b_{i}$ on the sets $\{\Theta \leq c\}$ and subspaces $E_{n}$,
(c) estimates $L_{n} V(x) \leq C-\kappa \Theta(x)$ for all $x \in E_{n}$ and $\left|b_{i}(x)\right| \leq C_{i}+\delta_{i}(\Theta(x)) \Theta(x)$ for $x \in\{\Theta<\infty\}$, where $\kappa>0$ and $\lim _{s \rightarrow+\infty} \delta_{i}(s)=0$. For all $x \in X$, let

$$
\begin{gathered}
V(x):=\int_{D} x(u)^{2} d u \\
\Theta(x):=\int_{D}\left|\nabla\left(\zeta_{(r+1) / 2} \circ x\right)(u)\right|^{2} d u
\end{gathered}
$$

where $\Theta(x):=+\infty$ if $\zeta_{(r+1) / 2} \circ x \notin H_{0}^{2,1}(D)$. By the Sobolev embedding theorem, the sets $\{\Theta \leq c\}$ are compact in $X$ and the functions $b_{i}$ are continuous on them (in the topology of $X$ ), hence also on $E_{n}$. Indeed, given a sequence of functions $x_{j} \in X$ such that $\zeta_{(r+1) / 2} \circ x_{j} \in H_{0}^{2,1}(D)$ and $\left\|\zeta_{(r+1) / 2} \circ x_{j}\right\|_{2,1}^{2} \leq c$, one can find a subsequence $\zeta_{(r+1) / 2} \circ x_{j_{k}}$ that converges in $L^{2}(D)$. Therefore, the sequence $x_{j_{k}}$ converges in $L^{2}(D)$ to some function $x$. Clearly, $\left\|\zeta_{(r+1) / 2} \circ x\right\|_{2,1}^{2} \leq c$. The continuity of $b_{i}$ on $\{\Theta \leq c\}$ is seen by the same reasoning. In addition, $V$ is a positive definite quadratic form on the spaces $E_{n}$. We have for all $x \in E_{n}$

$$
\sum_{i=1}^{n} x_{i} b_{i}(x)=\sum_{i=1}^{n}\left[\lambda_{i} x_{i}\left(\Psi \circ x, e_{i}\right)_{2}+x_{i}\left(\Phi \circ x, e_{i}\right)_{2}\right]=(\Delta x, \Psi \circ x)_{2}+(x, \Phi \circ x)_{2}
$$

Let us pick $\alpha>1$ and $\kappa>0$ such that $\alpha \delta+\kappa=4 \kappa_{0} \lambda_{1}(r+1)^{-2}$. One can find $C_{\alpha}>0$ such that $|x(u) \Phi(x(u))| \leq C_{\alpha}+\alpha \delta|x(u)|^{r+1}$. Taking into account the estimate

$$
\frac{(r+1)^{2}}{4} \int_{D}|x(u)|^{r-1}|\nabla x(u)|^{2} d u=\left.\left.\int_{D}|\nabla| x\right|^{(r+1) / 2}(u)\right|^{2} d u \geq \lambda_{1} \int_{D}|x(u)|^{r+1} d u
$$

$$
\begin{aligned}
L_{n} V(x) & =2 \sum_{i=1}^{n} q_{i}+2 \sum_{i=1}^{n} x_{i} b_{i}(x) \\
& =2 \sum_{i=1}^{n} q_{i}+2 \int_{D}[\Delta x(u) \Psi(x(u))+x(u) \Phi(x(u))] d u \\
& =2 \sum_{i=1}^{n} q_{i}-2 \int_{D} \Psi^{\prime}(x(u))|\nabla x(u)|^{2} d u+2 \int_{D} x(u) \Phi(x(u)) d u \\
& \leq 2 S-2 \kappa_{0} \int_{D}|x(u)|^{r-1}|\nabla x(u)|^{2} d u+2 C_{\alpha}|D|+2 \delta \alpha \int_{D}|x(u)|^{r+1} d u \\
& \leq \kappa^{\prime}-\kappa \Theta(x)
\end{aligned}
$$

where $\kappa^{\prime} \geq 0$ and $|D|$ is the measure of $D$. Finally, taking into account that $e_{i}$ is a bounded function, $\left|b_{i}(x)\right|$ can be estimated by $\alpha_{i}+\beta_{i} \int_{D}|x(u)|^{r} d u$ with some positive numbers $\alpha_{i}$ and $\beta_{i}$. It remains to observe that for all $x \in\{\Theta<\infty\}$ one has

$$
\int_{D}|x(u)|^{r} d u \leq\left(\int_{D}|x(u)|^{r+1} d u\right)^{r /(r+1)} \leq \lambda_{1}^{-r /(r+1)}\left(\left.\left.\int_{D}|\nabla| x\right|^{(r+1) / 2}(u)\right|^{2} d u\right)^{r /(r+1)}
$$

Therefore, we obtain

$$
\left|b_{i}(x)\right| \leq \alpha_{i}^{\prime}+\beta_{i}^{\prime} \Theta(x)^{-1 /(r+1)} \Theta(x), \quad x \in\{\Theta<\infty\}
$$

Thus, all the hypotheses of the theorem cited are satisfied. Therefore, we obtain a probability measure $\mu$ on $X$ satisfying equation (3) such that $\zeta_{(r+1) / 2} \circ x \in H_{0}^{2,1}(D)$ for $\mu$-a.e. $x$ and the function $\left\|\zeta_{(r+1) / 2} \circ x\right\|_{2,1}^{2}$ is $\mu$-integrable. In addition, $\mu$ is the weak limit of a subsequence of the sequence of probability measures $\mu_{n}$ on $E_{n}$ satisfying $L_{n}^{*} \mu_{n}=0$ and

$$
\begin{equation*}
\sup _{n} \int_{E_{n}} \int_{D}\left|\nabla\left(\zeta_{(r+1) / 2} \circ x\right)(u)\right|^{2} d u \mu_{n}(d x)=: K<\infty \tag{7}
\end{equation*}
$$

In particular,

$$
\begin{equation*}
\sup _{n} \int_{E_{n}} \int_{D}|x(u)|^{r-1} d u \mu_{n}(d x)=: K_{1}<\infty \tag{8}
\end{equation*}
$$

(ii) Now we suppose that $\sum_{i=1}^{\infty} q_{i} \sup _{u \in D} e_{i}(u)^{2}=: M<\infty$. First we consider the case where the function $|\Phi|$ is bounded. Let $\mu$ be the measure constructed in (i). We show that $\zeta_{r} \circ x \in H_{0}^{2,1}(D)$ for $\mu$-a.e. $x$ and the function $\left\|\zeta_{r} \circ x\right\|_{2,1}^{2}$ is $\mu$-integrable. We may assume that the whole sequence $\left\{\mu_{n}\right\}$ converges weakly to $\mu$. It is important that this sequence is uniformly tight on the space $L^{r}(D)$, hence converges weakly to $\mu$ also on that space. This is obvious from (7) and the Sobolev embedding theorem. Set

$$
\Xi(t):=\int_{0}^{t} \Psi(s) d s
$$

Let $n$ be fixed and let

$$
V_{n}(x):=\int_{D} \Xi(x(u)) d u, \quad x \in E_{n}
$$

Let $\Lambda:=\sqrt{-\Delta}$. We observe that for any $u \in H_{0}^{2,1}(D)$, one has

$$
\left(u, \Delta P_{n} u\right)_{2}=-\left(\Lambda u, \Lambda P_{n} u\right)_{2}=-\left(P_{n} \Lambda u, P_{n} \Lambda u\right)_{2}
$$

since $P_{n}=P_{n}^{2}$ and $P_{n}$ commutes with $\Lambda$. Therefore, for all $x \in E_{n}$ we have

$$
\begin{aligned}
& L_{n} V_{n}(x)=\sum_{i=1}^{n} q_{i} \int_{D} \Psi^{\prime}(x(u)) e_{i}(u)^{2} d u+\int_{D} \Psi(x(u)) \Delta P_{n}(\Psi \circ x)(u) d u \\
& +\left(P_{n}(\Phi \circ x), P_{n}(\Psi \circ x)\right)_{2} \\
& =\sum_{i=1}^{n} q_{i} \int_{D} \Psi^{\prime}(x(u)) e_{i}(u)^{2} d u-\left\|P_{n} \Lambda(\Psi \circ x)\right\|_{2}^{2}+\left(P_{n}(\Phi \circ x), P_{n}(\Psi \circ x)\right)_{2} \\
& \leq C_{0} M+\kappa_{1} M \int_{D}|x(u)|^{r-1} d u-\left\|P_{n} \Lambda(\Psi \circ x)\right\|_{2}^{2} \\
& +\frac{\lambda_{1}}{2}\left\|P_{n}(\Psi \circ x)\right\|_{2}^{2}+\frac{1}{2 \lambda_{1}}\|\Phi \circ x\|_{2}^{2} \\
& \leq C_{0} M+\kappa_{1} M \int_{D}|x(u)|^{r-1} d u-\frac{1}{2}\left\|P_{n} \Lambda(\Psi \circ x)\right\|_{2}^{2}+\frac{1}{2 \lambda_{1}}\|\Phi \circ x\|_{2}^{2} .
\end{aligned}
$$

It is easily seen (see, e.g., the proofs of Lemma 1.2 in [5] and Theorem 4.1 in [4]) that this estimate along with (8) yields

$$
\int_{E_{n}}\left\|P_{n} \Lambda(\Psi \circ x)\right\|_{2}^{2} \mu_{n}(d x) \leq 2 C_{0} M+2 \kappa_{1} M K_{1}+\frac{1}{\lambda_{1}} \int_{E_{n}}\|\Phi \circ x\|_{2}^{2} \mu_{n}(d x)
$$

Then for every $N$ we have

$$
\begin{equation*}
\int_{X}\left\|P_{N} \Lambda(\Psi \circ x)\right\|_{2}^{2} \mu(d x) \leq 2 C_{0} M+2 \kappa_{1} M K_{1}+\frac{1}{\lambda_{1}} \int_{X}\|\Phi \circ x\|_{2}^{2} \mu(d x) \tag{9}
\end{equation*}
$$

Indeed, if $n \geq N$, then $\left\|P_{N} \Lambda(\Psi \circ x)\right\|_{2}^{2} \leq\left\|P_{n} \Lambda(\Psi \circ x)\right\|_{2}^{2}$. Hence

$$
\int_{E_{n}}\left\|P_{N} \Lambda(\Psi \circ x)\right\|_{2}^{2} \mu_{n}(d x) \leq 2 C_{0} M+2 \kappa_{1} M K_{1}+\frac{1}{\lambda_{1}} \int_{E_{n}}\|\Phi \circ x\|_{2}^{2} \mu_{n}(d x) .
$$

We observe that

$$
g_{N}(x):=\left\|P_{N} \Lambda(\Psi \circ x)\right\|_{2}^{2}=\sum_{i=1}^{N} \lambda_{i}\left(\Psi \circ x, e_{i}\right)_{2}^{2}
$$

is a continuous function on $L^{r}(D)$. Hence by the weak convergence we arrive at (9). By Fatou's lemma we obtain $\|\Lambda(\Psi \circ x)\|_{2}^{2}<\infty$ for $\mu$-a.e. $x$ and

$$
\int_{X}\|\Lambda(\Psi \circ x)\|_{2}^{2} \mu(d x) \leq 2 C_{0} M+2 \kappa_{1} M K_{1}+\frac{1}{\lambda_{1}} \int_{X}\|\Phi \circ x\|_{2}^{2} \mu(d x) .
$$

By using the estimates $|\Phi(s)| \leq C+4 r \lambda_{1}(r+1)^{-2}|\Psi(s)|$ and $4 r(r+1)^{-2}<1$ for $r>1$ along with the inequality $\lambda_{1}\|\Psi \circ x\|_{2}^{2} \leq\|\Lambda(\Psi \circ x)\|_{2}^{2}$, we obtain

$$
\begin{equation*}
\int_{X}\|\Lambda(\Psi \circ x)\|_{2}^{2} \mu(d x) \leq N\left(r, C, C_{0}, M, \kappa_{1}, K_{1}, \lambda_{1}\right) \tag{10}
\end{equation*}
$$

In the case where $\Phi$ is not bounded, we apply the above proved assertion to the functions $\Phi_{j}$ defined as follows: $\Phi_{j}(t)=\Phi(t)$ if $|\Phi(t)| \leq j, \Phi_{j}(t)=j \operatorname{sgn} \Phi(t)$ if $|\Phi(t)|>j$. Due to (10) the obtained measures $\mu_{j}$ form a uniformly tight sequence. We take for $\mu$ a limit point of $\left\{\mu_{j}\right\}$ in the weak topology. It is clear that for the new measures $\mu_{j}$ one has the uniform estimate (7), which yields (4) (we do not claim that this measure $\mu$ coincides with the measure constructed in (i)). Estimate (6) follows by Lemma 1.

For example, if $\Psi(t)=t^{m}$, where $m$ is an odd number, then one can take for $\Phi$ any polynomial of degree $m$ with a sufficiently small leading coefficient (the smallness of which depends on $\lambda_{1}$, in particular, one can take $\Phi(x)=x^{m}$ provided $\lambda_{1}$ is sufficiently large).

Now we can pass from cylindrical functions to $C_{b}^{2}$ functions, however, defined on larger spaces such as $H$. Note that by Lemma 1 and assertion (ii) of Theorem 1 we have $\Delta \Psi \circ x \in H$ for $\mu$-a.e. $x$. By finite dimensional approximations we obtain the following.

Corollary 1. In the situation of assertion (ii) of Theorem 1, the constructed measure $\mu$ satisfies the equation $L^{*} \mu=0$ on $H$ with respect to the class $C_{b}^{2}(H)$.

Corollary 2. Suppose that in either assertion of Theorem 1 one has $\Psi^{\prime}(0)>0$. Then $\mu\left(H_{0}^{2,1}(D)\right)=1$ and

$$
\int_{X} \int_{D}|\nabla x(u)|^{2} d u \mu(d x)<\infty
$$

Proof. Due to our assumption, in a neighborhood of zero $\Psi^{\prime}(s) \geq \alpha$ for some constant $\alpha>0$. Now the assertion follows by the same reasoning as in the theorem with the function

$$
\Theta(x)=\int_{D}\left[|\nabla x(u)|^{2}+\left|\nabla\left(\zeta_{(r+1) / 2} \circ x\right)(u)\right|^{2}\right] d u
$$

It is very likely that as in [1], [2], for any $k \in \mathbb{N}, k \geq 2$, one has $x^{k} \in H_{0}^{2,1}(D)$ for $\mu$-a.e. $x$ and $\left\|x^{k}\right\|_{2,1}^{2}$ is $\mu$-integrable. If we formally consider the Lyapunov functions $\int_{D} x^{2 k}(t) d t$, then we obtain for $\Theta$ the function $\int_{D}|x(u)|^{r+2 k-2}|\nabla x(u)|^{2} d u$. However, we have not managed to justify this (except for the case $k=r$ in assertion (ii) above).

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