

# About substitution tilings with statistical circular symmetry

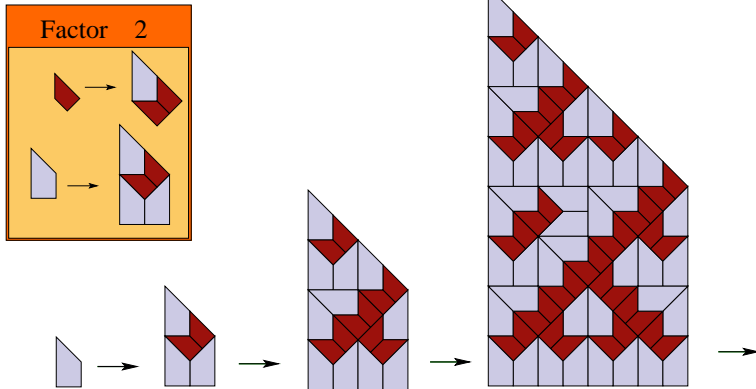
Dirk Frettlöh

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Bielefeld, Germany

Quasicrystals - The Silver Jubilee  
Tel Aviv  
14-19 October 2007

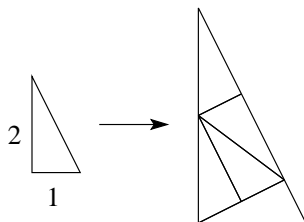
1. Examples of tilings with statistical circular symmetry
2. Diffraction of...
3. (Dynamics of... )

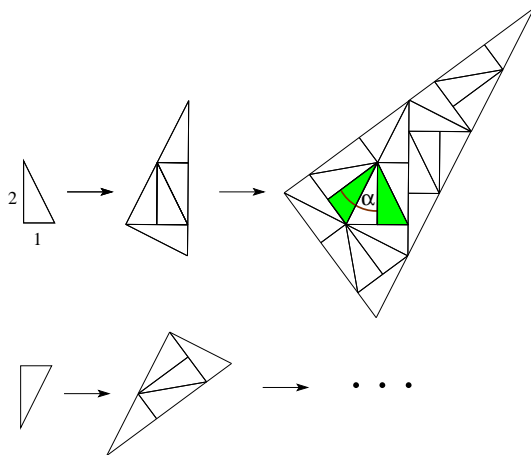
## Substitution tilings:



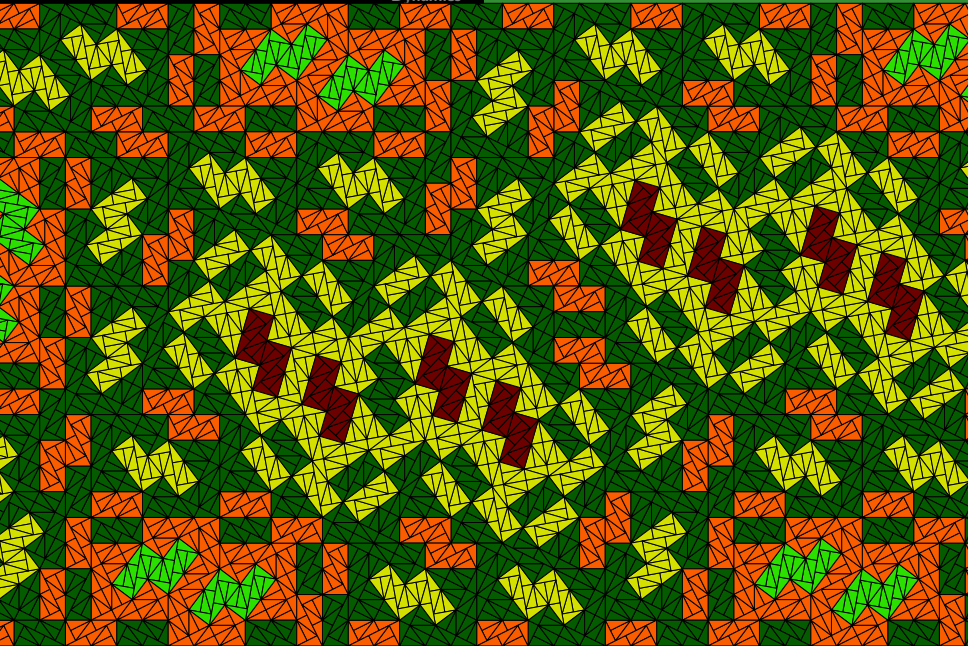
Usually, tiles occur in finitely many different orientations.

Not always. Conway's Pinwheel substitution (1991):

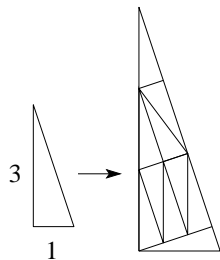




The angle  $\alpha$  is  
*irrational*; that is,  
 $\alpha \notin \pi\mathbb{Q}$ .

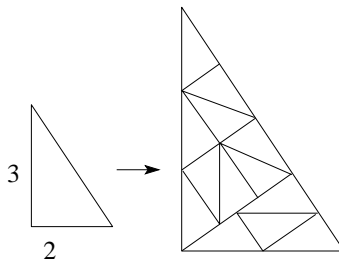


Obvious generalizations: Pinwheel  $(n, k)$



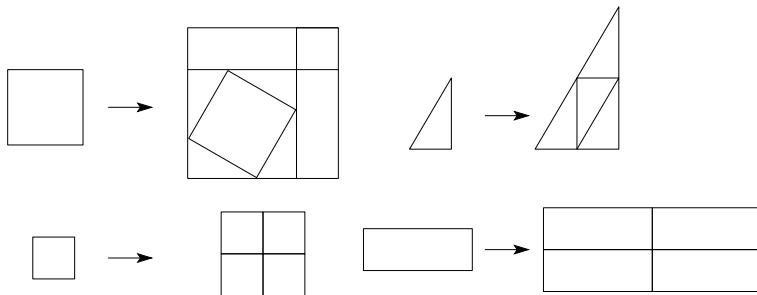
$n = 3, k = 1$

etc.



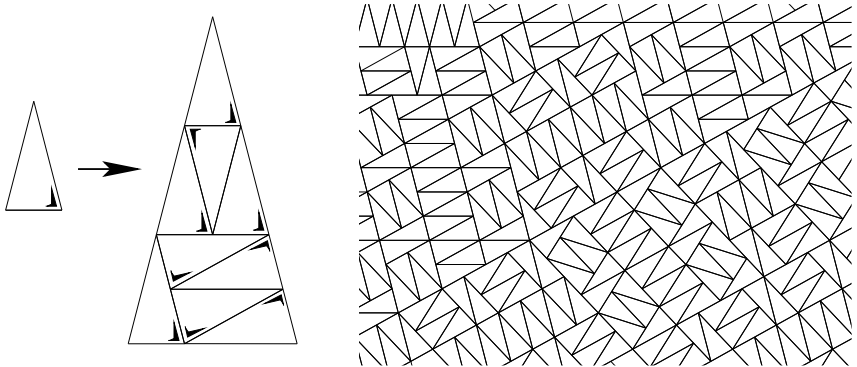
$n = 3, k = 2$

Cesi's example (1990):



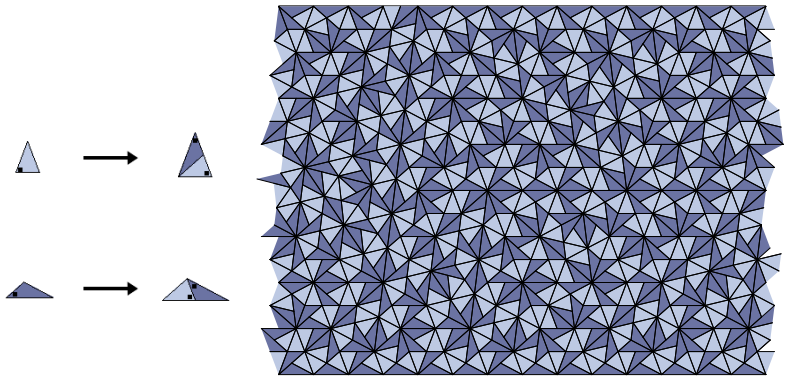


Unknown (< 1996, communicated to me by Danzer):

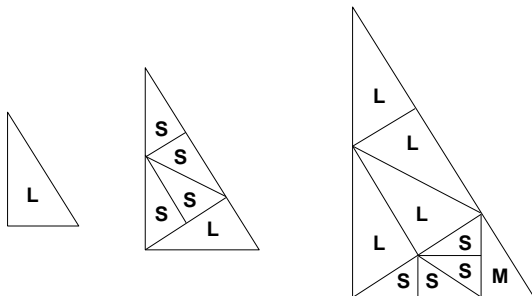


(+ obvious generalizations)

C. Goodman-Strauss, L. Danzer (ca. 1996):

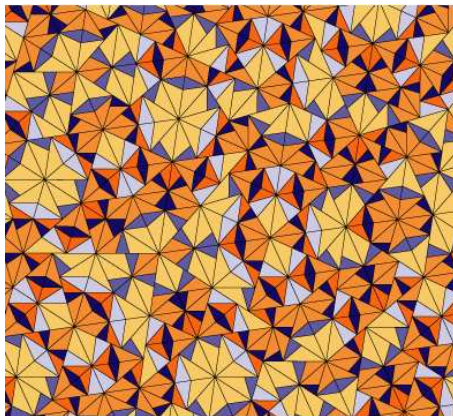
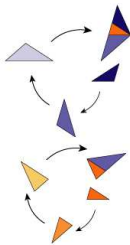


Sadun's generalized Pinwheels (1998):

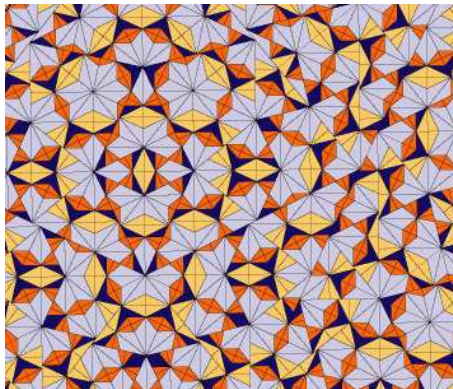
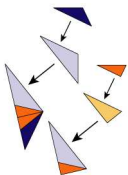


Yields infinitely many proper tile-substitutions.

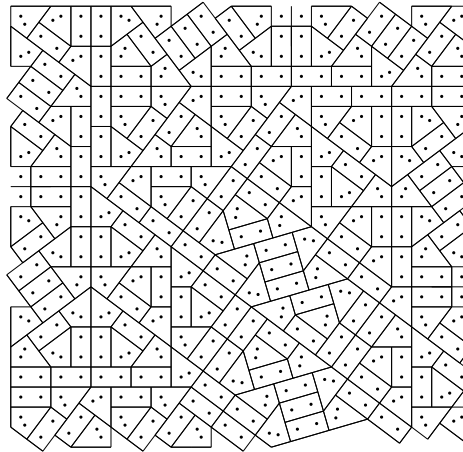
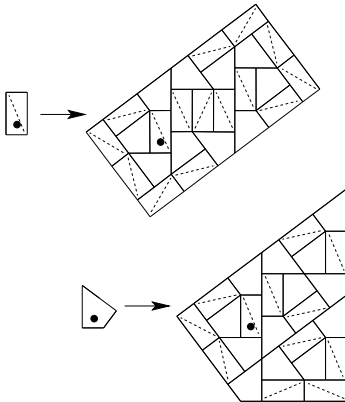
# Harriss' Cubic Pinwheel (2004 $\pm 1$ ):



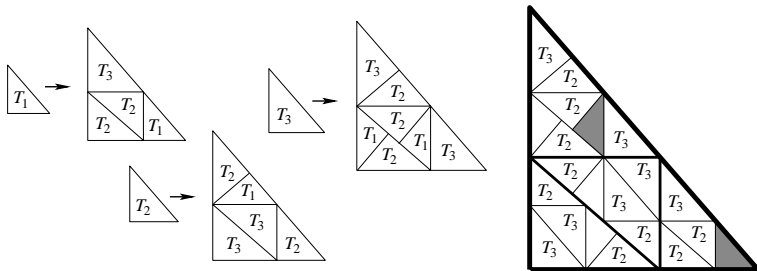
# Harriss' Quartic Pinwheel (2004 $\pm 1$ ):



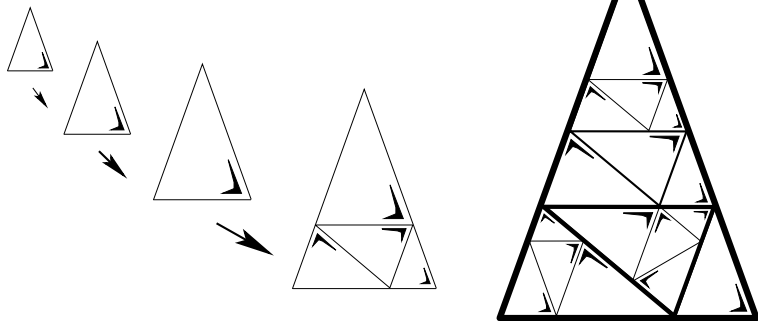
Kite domino (equivalent with pinwheel):



Pythia  $(m,j)$ , here:  $m = 3, j = 1$ .

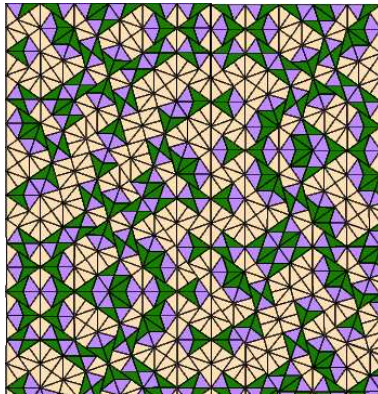
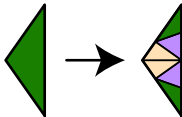
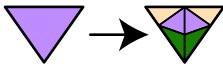
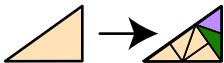


Tipi  $(m,j)$ , here:  $m = 3, j = 1$ .

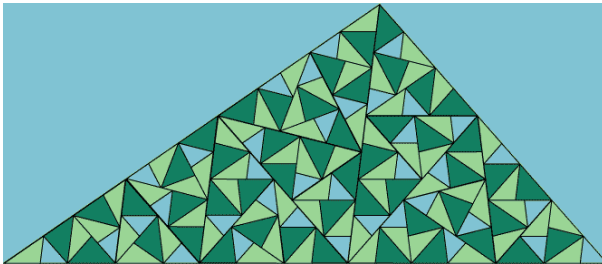




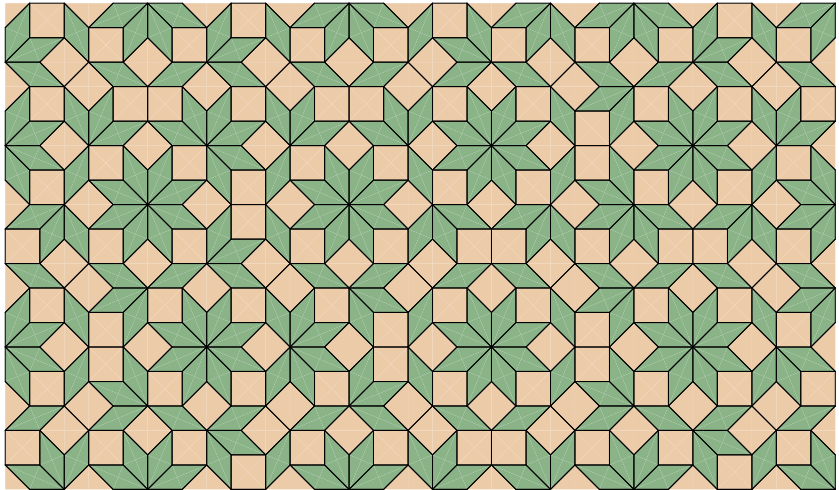
## Dale Walton: several single examples

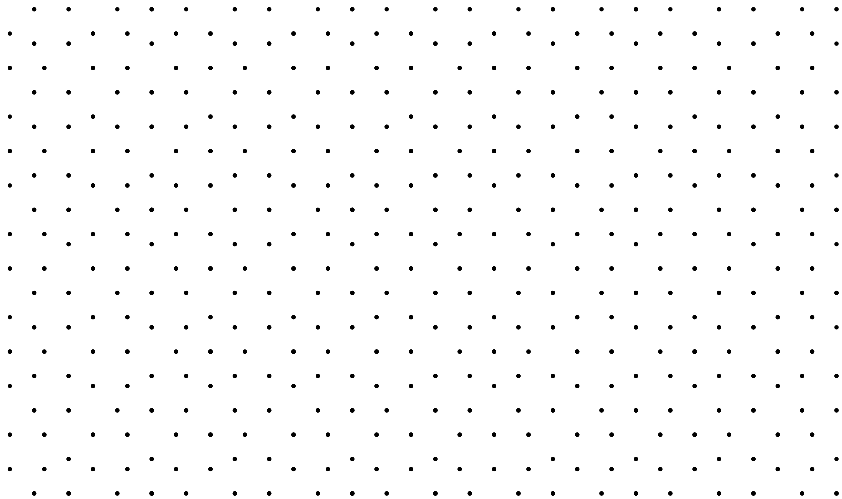


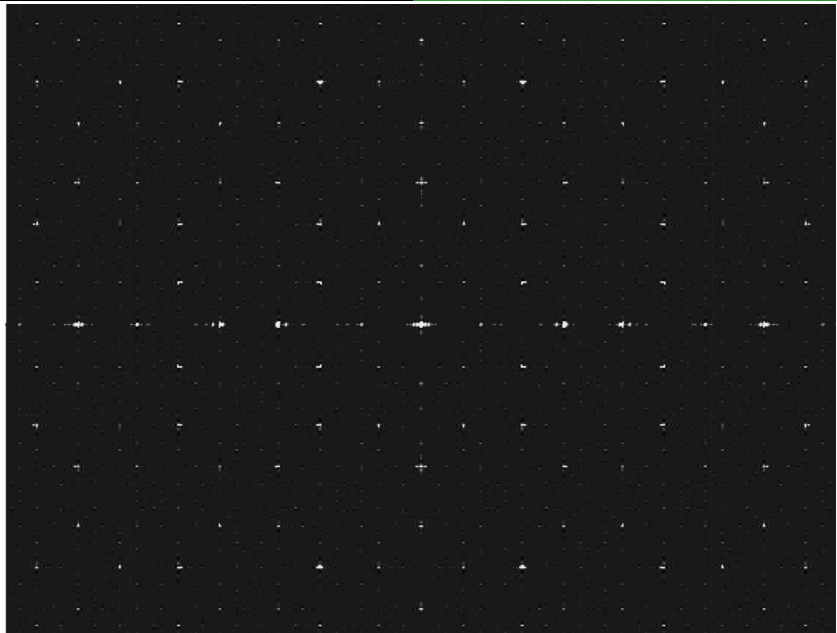
Dale Walton: several single examples



## 2. Diffraction







## Mathematical description:

- ▶ Tiling  $\leadsto$  discrete point set  $\Lambda$ .
- ▶ Autocorrelation  $\gamma_\Lambda = \lim_{r \rightarrow \infty} \frac{1}{\text{vol } B_r} \sum_{x,y \in \Lambda \cap B_r} \delta_{x-y}$ .
- ▶ Fouriertransform  $\hat{\gamma}_\Lambda$  of the autocorrelation is the *diffraction spectrum*.

Mathematical description:

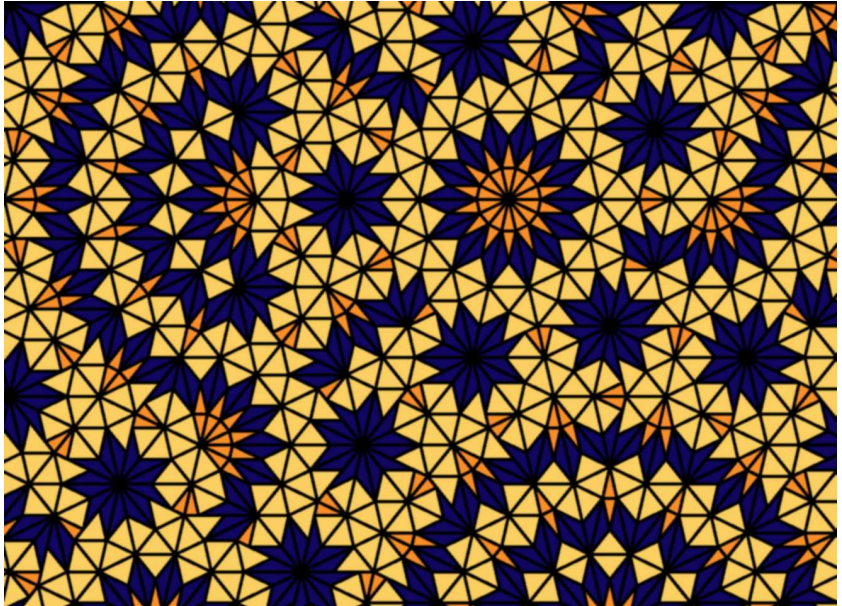
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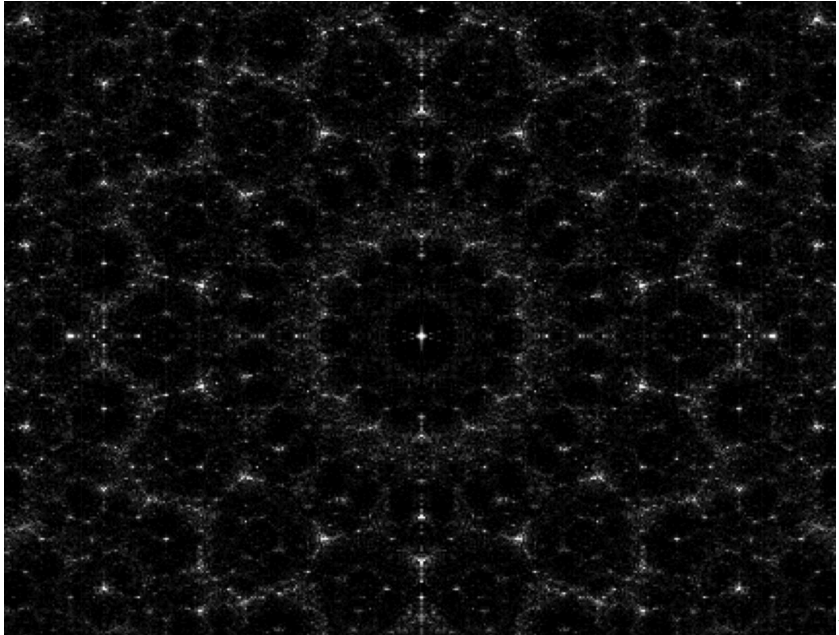
Since  $\hat{\gamma} := \hat{\gamma}_\Lambda$  is again a measure, it decomposes into three parts:

$$\hat{\gamma} = \hat{\gamma}_{pp} + \hat{\gamma}_{sc} + \hat{\gamma}_{ac}$$

(pp: pure point, ac: absolutely continuous, sc: singular continuous)







$$\hat{\gamma} = \hat{\gamma}_{pp} + \hat{\gamma}_{sc} + \hat{\gamma}_{ac}$$

For an ideal (mathematical, infinite) quasicrystal:

$$\hat{\gamma} = \hat{\gamma}_{pp}$$

For primitive substitution tilings with infinitely many orientations (we've just seen many examples):

$$\hat{\gamma} = \delta_0 + \hat{\gamma}_{sc} + \hat{\gamma}_{ac},$$

and  $\hat{\gamma}$  is circular symmetric.

This follows from:

### Theorem

*Each primitive substitution tiling with infinitely many orientations is of **statistical circular symmetry**, i.e. the orientations are equidistributed on the circle.*

Roughly spoken, each orientation occurs with the same probability.

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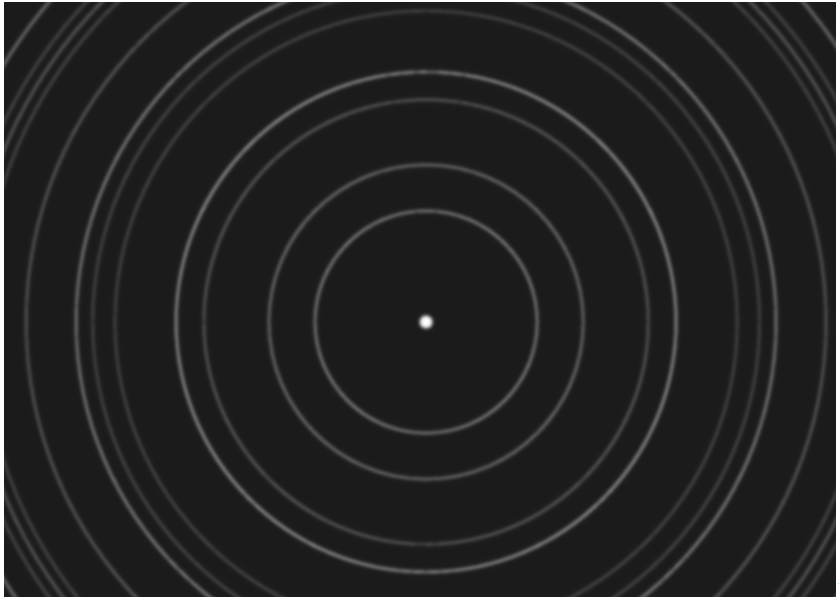
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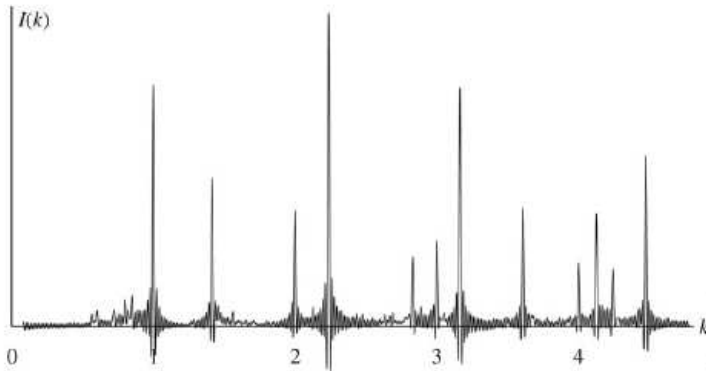
Consequently, the autocorrelation is of perfect circular symmetry, and circular symmetry of the diffraction spectrum follows.

Continuous parts are still mysterious.

## Pinwheel diffraction



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### 3. Dynamics / more about symmetry

The *hull*  $\mathbb{X}_{\mathcal{T}}$  of a tiling:

The closure of  $\{\mathcal{T} + t \mid t \in \mathbb{R}^d\}$  in an appropriate topology.

For primitive substitution tilings:

hull = LI-class = set of all tilings given by this substitution.

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In the pure point case, the dynamical spectrum of the dynamical system  $(\mathbb{X}_{\mathcal{T}}, \mathbb{R}^d)$  yields information about the diffraction spectrum of  $\mathcal{T}$ . (Dworkin, Hof, Solomyak, Robinson, Schlottmann...).

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Maybe we can extend this to the continuous parts of the spectrum.