Hunting French Ducks in a Noisy Environment

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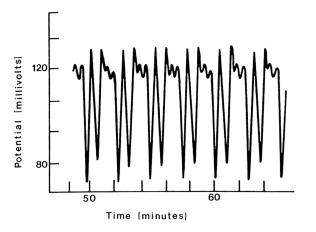
joint work with: Nils Berglund, MAPMO-CNRS, Orléans Barbara Gentz, Universität Bielefeld

Outline

- 1. Motivation: Mixed-Mode Oscillations
- 2. Introduction to Multiple Time Scale Dynamics
- 3. Canards near a Folded Node
- 4. Stochastic Blow-Up and Linearization
- 5. Covariance Tubes
- 6. (Early Jumps)

Mixed-Mode Oscillations (MMOs)

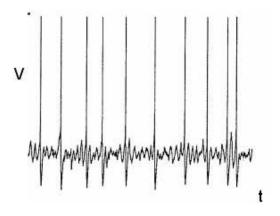
Belousov-Zhabotinsky reaction (Hudson, Hart and Marinko 1979):



Notation: $\ldots L_{j-1}^{s_{j-1}} L_j^{s_j} L_{j+1}^{s_{j+1}} \ldots$ here: $L^s = 2^2$.

Mixed-Mode Oscillations (MMOs)

Layer II Stellate Cells (Dickson et al. 2000):



Q: What is the mechanism for the small-amplitude oscillations?



Fast-Slow Systems

A general **fast-slow system** is a special ODE:

$$\begin{array}{rcl} \frac{dx}{dt} & = & x' & = & f(x,y) \\ \frac{dy}{dt} & = & y' & = & \epsilon g(x,y) \end{array}$$

where $(x,y) \in \mathbb{R}^m \times \mathbb{R}^n$ and $0 < \epsilon \ll 1$.

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where $(x,y) \in \mathbb{R}^m \times \mathbb{R}^n$ and $0 < \epsilon \ll 1$.

On the **slow time scale** $s = \epsilon t$ we get:

$$\begin{array}{rcl} \epsilon \frac{dx}{ds} & = & \epsilon \dot{x} & = & f(x,y) \\ \frac{dy}{ds} & = & \dot{y} & = & g(x,y) \end{array}$$

The Singular Limits

$$\overset{\epsilon=0}{\Rightarrow}$$

fast subsystem

$$x' = f(x, y)$$
$$y' = 0$$

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$$x' = f(x,y)$$

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$$\epsilon \dot{x} = f(x, y)$$

 $\dot{y} = g(x, y)$

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slow subsystem

$$0 = f(x, y)$$

$$\dot{y} = \sigma(x, y)$$

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slow subsystem

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$$\dot{y} = g(x, y)$$

Idea: Combine the two systems to analyze the case $0 < \epsilon \ll 1$.

Define the critical manifold

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 C_0 is **normally hyperbolic** at $P \in C_0$ if

 $(D_x f)(P)$ has no eigenvalues λ_j with zero real parts.

- C_0 is **attracting** if $\lambda_j < 0$ for all j.
- C_0 is **repelling** if there exists $\lambda_j > 0$.

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Theorem (Fenichel's Theorem, 1979)

A normally hyperbolic critical manifold C_0 perturbs $(0 < \epsilon \ll 1)$ to a **slow manifold** C_{ϵ} . C_{ϵ} is an $O(\epsilon)$ -distance away from C_0 and the slow subsystem flow approximates the flow on C_{ϵ} .

An Example - The Planar Fold

$$\begin{array}{rcl}
\epsilon \dot{x} & = & y - x^2 \\
\dot{y} & = & \mu - x
\end{array}$$

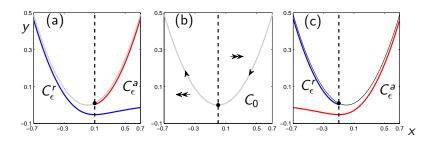


Figure: $\epsilon = 0.05$. (a) $\mu = 0.1$ (b) $\mu = 0$ (c) $\mu = -0.1$.

Folded Singularties in \mathbb{R}^3

Consider the following normal form:

$$\begin{array}{rcl} \epsilon \dot{x} & = & y-x^2, \\ \dot{y} & = & -(\mu+1)x-z, \\ \dot{z} & = & \frac{\mu}{2}, \end{array}$$

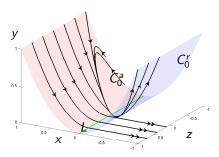
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The critical manifold decomposes as:

$$C_0 = \{(x, y, z) \in \mathbb{R}^3 : y = x^2\} = C_0^a \cup L \cup C_0^r$$



Let's calculate the slow flow

$$0 = y - x^2, \qquad \Rightarrow \quad \dot{y} = 2x\dot{x}.$$

Therefore the slow subsystem is

$$2x\dot{x} = -(\mu+1)x - z,$$

$$\dot{z} = \frac{\mu}{2}.$$

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Set $s \to 2x$ s; the **desingularized slow subsystem** is

$$\dot{x} = -(\mu + 1)x - z,
\dot{z} = \mu x.$$

Equilibrium (x, z) = (0, 0) for desingularized slow flow. Eigenvalues are

$$(\lambda_s, \lambda_w) := (-1, -\mu).$$

The origin (0,0) is a **folded node** for $\mu \in (0,1)$.

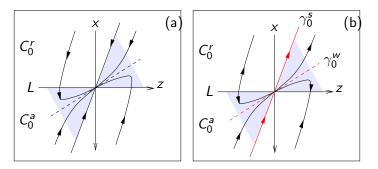


Figure: Strong singular canard γ_0^s ; weak singular canard γ_0^w .

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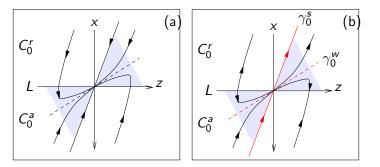


Figure: Strong singular canard γ_0^s ; weak singular canard γ_0^w .

Definition: A **maximal canard** is an orbit in $C^a_{\epsilon} \cap C^r_{\epsilon}$.



Theorem (Benoît 1990; Szmolyan/Krupa/Wechselberger 2000)

For $\epsilon>0$ sufficiently small the singular strong canards $\gamma_0^{s,w}$ perturb to maximal canards $\gamma_\epsilon^{s,w}$. Suppose $k\in\mathbb{N}$ and

$$2k+1 < \mu^{-1} < 2k+3$$
 and $\mu^{-1} \neq 2(k+1)$.

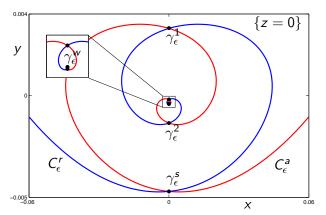
there are k other maximal canards that rotate around $\gamma_{\epsilon}^{\mathrm{W}}$.

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Geometric Desingularization (or Blow-Up)

Recall the normal form

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and apply

$$(x, y, z, s) = (\sqrt{\epsilon}\bar{x}, \epsilon\bar{y}, \sqrt{\epsilon}\bar{z}, \sqrt{\epsilon}\bar{s})$$

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$$(x, y, z, s) = (\sqrt{\epsilon}\bar{x}, \epsilon\bar{y}, \sqrt{\epsilon}\bar{z}, \sqrt{\epsilon}\bar{s})$$

This yields (dropping overbars for convenience)

$$\begin{array}{rcl} \dot{x} & = & y - x^2, \\ \dot{y} & = & -(\mu + 1)x - z, \\ \dot{z} & = & \frac{\mu}{2}. \end{array}$$

Q: Spacing of canards on $\{z = 0\}$?

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The linearized variational equation around the weak canard γ_0^w is

$$\mu \frac{du}{dz} = \underbrace{\begin{pmatrix} 4z & 2 \\ -2(\mu+1) & 0 \end{pmatrix}}_{=:A(z)} u = A(z)u.$$

Eigenvalues are $2z \pm i\omega(z) \Rightarrow$ contraction (z < 0) + rotation!?

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Lemma (Canard Spacing)

On $\{z=0\}$ the distance of the k-th maximal canard to γ_0^w is

$$\mathcal{O}(e^{-c_0(2k+1)^2\mu})$$

Stochastic Folded Nodes

Consider the normal form

$$dx_s = \frac{1}{\epsilon}(y_s - x_s^2)ds + \frac{\sigma}{\sqrt{\epsilon}}dW_s^{(1)},$$

$$dy_s = [-(\mu + 1)x_s - z_s]ds + \sigma'dW_s^{(2)},$$

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Main Idea: Control sample paths near deterministic solution.

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Main Idea: Control sample paths near deterministic solution.

Strategy:

- 1. Geometric desingularization (Blow-Up).
- Linearization around deterministic solution.
- 3. Covariance evolution provides tubular neighbourhoods.
- 4. Stay inside tubes for $-1 < z < \sqrt{\mu}$.
- 5. Need to control nonlinearity and diffusion.

Blow-up (rescale) the normal form as before

$$(x, y, z, s) = (\sqrt{\epsilon}\bar{x}, \epsilon\bar{y}, \sqrt{\epsilon}\bar{z}, \sqrt{\epsilon}\bar{s})$$

then (dropping overbars for convenience)

$$dx_{s} = (y_{s} - x_{s}^{2})ds + \frac{\sigma}{\epsilon^{3/4}}dW_{s}^{(1)},$$

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We also re-scale the noise level parameters and set

$$(\epsilon^{3/4}\sigma, \epsilon^{3/4}\sigma') =: (\bar{\sigma}, \bar{\sigma}')$$

Observe: Can use s or z as "time" variable.

The Stochastic Variational Equation

Focusing on
$$(x_z, y_z) = (x_z^{\text{det}} + \xi_z, y_z^{\text{det}} + \eta_z)$$
 we get
$$d\xi_z = \frac{2}{\mu} (\eta_z - \xi_z^2 - 2x_z^{\text{det}} \xi_z) dz + \frac{\sqrt{2}\sigma}{\sqrt{\mu}} dW_z^{(1)},$$

$$d\eta_z = -\frac{2}{\mu} (\mu + 1) \xi_z dz + \frac{\sqrt{2}\sigma'}{\sqrt{\mu}} dW_z^{(2)}.$$

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Proposition

Linearize the variational equation; set $V(z) := \sigma^{-2} Cov(z)$ then

$$\begin{array}{lcl} \dot{v}_{11} & = & -8x^{\mathsf{det}}(z)v_{11} + 4v_{12} + 2, \\ \dot{v}_{22} & = & -4(\mu+1)v_{12} + 2(\sigma'/\sigma)^2, \\ \dot{v}_{12} & = & -2(\mu+1)v_{11} + 2v_{22} - 4x^{\mathsf{det}}(z)v_{12}. \end{array}$$

Note
$$v_{12} = Cov(\xi_z, \eta_z) = Cov(\eta_z, \xi_z) = v_{21}$$
.

Neighbourhoods of Deterministic Solutions

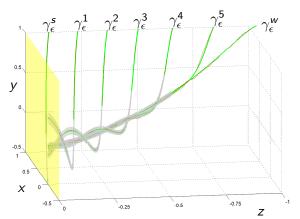
Let (x(z), y(z)) =: w(z) be a deterministic solution. Define a **tube-shaped**-neighbourhood

$$\mathcal{B}(r) = \{(x, y, z) : z_0 \le z \le \sqrt{\mu}, \\ [(x, y) - w(z)] \cdot V(z)^{-1} [(x, y) - w(z)] < r^2 \}.$$

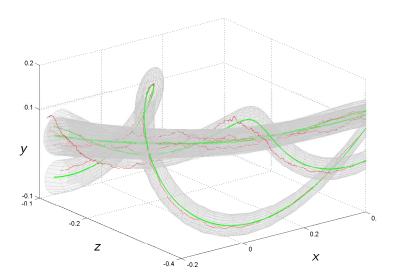
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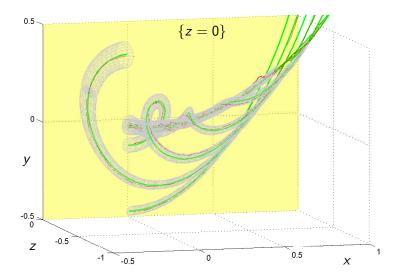
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Zoom



Front View



Covariance Estimates

Theorem (Covariance Tubes)

On the section $\{z=0\}=\{\bar{z}=0\}$ we have (as $\mu \to 0$):

$$v_1=\mathcal{O}(1/\sqrt{\mu}), \quad v_2=\mathcal{O}(1/\sqrt{\mu}), \quad v_3=\mathcal{O}(1), \quad (v_1-v_2)=\mathcal{O}(1)$$

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Sketch of Proof.

- 1. Change coordinates in the variational equation.
- 2. Use a symmetry to reduce it to a planar system.
- 3. A complex eigenvalue pair crosses the imaginary axis at z = 0.
- 4. View the planar system as a fast subsystem with slow time z.
- 5. Apply the delayed Hopf bifurcation theory.

The Nonlinear Variational SDE

It turns out that in suitable coordinates we have to deal with

$$d\zeta_z = \frac{1}{\mu} \big[A(z)\zeta_z + b(\zeta_z, z) \big] dz + \frac{\sigma}{\sqrt{\mu}} F(z) dW_z,$$

where $\zeta_z = (\xi_z, \eta_z)$ and A(z) is now given by

$$A(z) = \begin{pmatrix} -2x(z) & \omega_2(z) \\ -\omega_2(z) & -2x(z) \end{pmatrix}.$$

Staying inside $\mathcal{B}(r)$...

Theorem (Staying inside Covariance Tubes)

There exists a function $K(z,z_0)$ such that for $\kappa=1-\mathcal{O}(\cdot)$

$$\mathbb{P}\left\{\tau_{\mathcal{B}(r)} < z\right\} \leq K(z, z_0) \exp\left\{-\kappa \frac{r^2}{2\sigma^2}\right\}$$

holds for all z such that $z_0 \le z \le \sqrt{\mu}$.

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Sketch of Proof.

- 1. Consider a short time interval $[z_1, z_2]$.
- 2. Consider the fundamental solution U(z, u) for $\mu \dot{\zeta} = A(z)\zeta$.
- 3. Set $\Upsilon_u := U(z, u)\zeta_u$ and observe $\Upsilon_u = \Upsilon_u^0 + \Upsilon_u^1$

$$\Upsilon_u^0 = \frac{\sigma}{\sqrt{\mu}} \int_{z_0}^u U(z, v) F(v) dW_v,$$

$$\Upsilon_u^1 = \frac{1}{\mu} \int_{z_0}^u U(z, v) b(\zeta_v, v) dv.$$



Sketch of Proof (continued).

- 4. Υ_{μ}^{0} is a Gaussian martingale.
- 5. Doob's submartingale inequality, let $M_u := \|Q(z_1,z_2)\Upsilon_u^0\|$

$$\mathbb{P}\left\{\sup_{z_1 \leq u \leq z_2} e^{M_u^2} \geq e^{r^2}\right\} \leq \frac{1}{e^{r^2}} \mathbb{E}\left[e^{M_{z_2}}\right] \leq (\cdots) \mathcal{O}\left(e^{-\frac{r^2}{2\sigma^2}}\right)$$

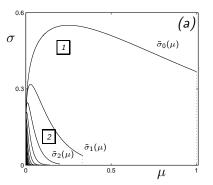
where $Q(z_1, z_2)$ is defined via the covariance matrix V.

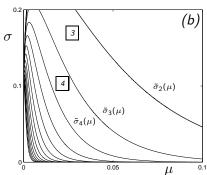
- 6. From last step we bound $\mathbb{P}\left\{\sup_{z_1 \leq u \leq z_2} M_u \geq r\right\}$.
- 7. Estimate $||Q(z_1, z_2)\Upsilon_u^1||$ directly and show that it is small.
- 8. We find that escape during a short time is highly unlikely.
- 9. Piece previous result together for a "nice" partition of $[z_0, z]$.



Theorem (Noise, Canards and SAOs)

Depending on noise intensity $\tilde{\sigma}$ and bifurcation parameter μ the "noisy interactions" of canards are:





$$\tilde{\sigma}_k(\mu) = \mu^{1/4} e^{-(2k+1)^2 \mu}$$

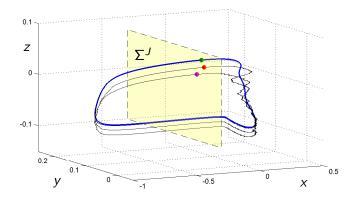
Further Result - Early Jumps

For $z > \sqrt{\mu}$ beyond the folded node, SDE paths jump early.

$$dx = \frac{1}{\epsilon}(y - x^2 - x^3)ds + \frac{\sigma}{\sqrt{\epsilon}}dW_s^{(1)},$$

$$dy = [-(\mu + 1)x - z]ds + \sigma'dW_s^{(2)},$$

$$dz = [\frac{\mu}{2} + ax + bx^2]ds.$$



Theorem (Escape of Sample Paths)

In blow-up coordinates, consider $z>\sqrt{\mu}$ and let $\mathcal D$ be a tube around γ^w that grows like $\mathcal O(\sqrt{z})$. Then the probability that a sample path stays in $\mathcal D$ becomes small as soon as

$$z \gg \sqrt{\mu |\log \sigma|/\nu}$$
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where $\nu > 0$.

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where $\nu > 0$.

Sketch of Proof.

- 1a. **Diffusion-dominated escape** from small set near γ^w .
- 1b. Subdivide again, need Markov property to re-start.
- 2a. **Drift-dominated escape** from \mathcal{D} .
- 2b. Change to polar coordinates.
- 2c. Use averaging to consider radius SDE.
- 2d. Show that drift dominates diffusion.



Back to Mixed-Mode Oscillations...

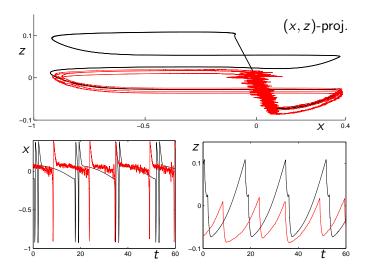


Figure: 3D model system again, different parameters ...

Conclusions

Overview I:

- ► Fast-slow systems can have intricate singularities.
- ▶ The SAOs of MMOs are often caused by these mechanisms.
- ▶ Deterministic scenario is often unrealistic (biophysics!).

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Overview I:

- Fast-slow systems can have intricate singularities.
- ▶ The SAOs of MMOs are often caused by these mechanisms.
- Deterministic scenario is often unrealistic (biophysics!).

Overview II:

- Metastable sample paths for SDEs are natural extension.
- Variational equations around solutions play a key role.
- Use Doob's inequality to control sample paths.
- Early jumps after passage through folded node region.
- ▶ Intricate dependencies between σ , μ and ϵ .

Main Reference:

(1) N. Berglund, B. Gentz, C. Kuehn, *Hunting French ducks in a noisy environment*, in preparation.

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Thank you for your attention.