



## Amplitude Equation for the generalized Swift Hohenberg Equation with Noise

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# Introduction

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The generalized Swift Hohenberg equation

$$\partial_t u = ru - (1 + \partial_x^2)^2 u + \alpha u^2 - u^3 \quad (\text{SH})$$

is a qualitative model for Rayleigh Benard convection.



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It is well known (Cross/Hohenberg 93, Hilali 95, Burke/Knobloch 06) that

$$u(t, x) \approx \sqrt{|r|} \cdot A(|r|t) \cdot e^{ix} + \sqrt{|r|} \cdot \overline{A(|r|t)} \cdot e^{-ix}.$$

where the complex amplitude  $A(T)$  of the dominant mode  $e^{ix}$  is the solution of

$$\partial_T A = \text{sign}(r)A + 3\left(\frac{38}{27}\alpha^2 - 1\right)|A|^2 A,$$



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We consider the following stochastic version of (SH)

$$\partial_t u = \nu \varepsilon^2 u - (1 + \partial_x^2)^2 u + \alpha u^2 - u^3 + \varepsilon \sigma \partial_t \beta, \quad (\text{SSH})$$

where

- ▶  $\beta(t)$  is a real valued standard Brownian motion,
- ▶  $\alpha$ ,  $\sigma$  and  $\nu$  are real-valued constants,
- ▶ the small parameter  $\varepsilon > 0$  relates the distance from bifurcation to the noise strength.

We estimate (SSH) by a similar amplitude equation as in the deterministic case:

$$dA = \left( \nu A + 3\left(\frac{38}{27}\alpha^2 - 1\right)A|A|^2 + 3\left(\alpha^2 - \frac{1}{2}\right)\sigma^2 A \right) dT + 2\alpha\sigma A d\tilde{\beta}. \quad (\text{AE})$$



# Result on bounded domains

## Theorem

Let  $T_0 > 0$ ,  $\alpha \in \mathbb{R}$  with  $\alpha^2 < \frac{27}{38}$  and  $0 < \kappa$ . Let  $u$  be a mild solution of (SSH) with  $\|u(0)\|_\infty = \mathcal{O}(\varepsilon^{1-\kappa})$ . Let  $A(T) \in \mathbb{C}$ ,  $T \in [0, T_0]$  solve (AE), then  $\forall p \in \mathbb{N} : \exists C_p$  such that

$$\mathbb{P} \left( \sup_{t \in [0, T_0]} \|u(t) - u_A(t) - e^{-t(1+\partial_x^2)^2} u_s(0)\|_\infty > \varepsilon^{2-19\kappa} \right) \leq C_p \varepsilon^p$$

with the approximation

$$u_A(t, x) = \varepsilon A(\varepsilon^2 t) e^{ix} + \varepsilon \bar{A}(\varepsilon^2 t) e^{-ix} + \varepsilon Z_\varepsilon(\varepsilon^2 t)$$

where  $Z_\varepsilon$  is the Ornstein-Uhlenbeck process defined by

$$Z_\varepsilon(T) := \varepsilon^{-1} \sigma \int_0^T e^{-\varepsilon^{-2}(T-s)} d\tilde{\beta}(s).$$

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# Comparing the two Amplitude equations

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Deterministic:

$$\partial_T A = rA + 3\left(\frac{38}{27}\alpha^2 - 1\right)|A|^2 A$$

With added noise:

$$dA = \left(\nu A + 3\left(\frac{38}{27}\alpha^2 - 1\right)A|A|^2\right)dT + 3\left(\alpha^2 - \frac{1}{2}\right)\sigma^2 A dT + 2\alpha\sigma A d\tilde{\beta}.$$

Where does the difference come from?

- ▶ Noise - Nonlinearity interaction,
- ▶ Averaging ( $\int a Z_\varepsilon^2 dt \approx \int a \sigma dt$ ).



# Idea of Proof / Rescaling

We rescale the the solutions of (SSH) to their natural timescale:

$$v(T) := \varepsilon^{-1} u(\varepsilon^2 T).$$

Since we are on bounded domains we can write  $v$  as

$$v = ae^{ix} + \varepsilon \Phi e^{i2x} + c.c. + \varepsilon \Psi + Z_\varepsilon + \sum_{|k| \geq 3} v_k e^{ikx} + e^{-T\varepsilon^{-2}(1+\partial_x^2)^2} v_s(0)$$

The mild solution of  $v_k$  looks as follows

$$v_k(T) = \int_0^T e^{-\varepsilon^{-2}(1-k^2)^2(T-s)} \left[ \nu v_k(s) + \varepsilon^{-1} \alpha(\widehat{v^2})_k(s) - (\widehat{v^3})_k(s) \right] ds,$$

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# Idea of Proof / Reduction

With this we can show that until a stopping time  $\tau_*$  ( $\|v(T)\|_\infty = \mathcal{O}(\varepsilon^{0-})$  for  $T \in [0, \tau_*]$ ) we have

$$\|v - a - Z - e^{-T\varepsilon^{-2}(1+\partial_x^2)^2} v_s(0)\|_\infty = \mathcal{O}(\varepsilon^{1-})$$

and by calculating  $(\widehat{v}^2)_i$  and  $(\widehat{v}^3)_i$  ( $i \in 1, 2, 3$ ) we get

$$da = (\nu a + 2\alpha \bar{a} \Phi + 2\alpha a \Psi - 3|a|^2 - 3aZ_\varepsilon^2 + \varepsilon^{-1}2\alpha a Z_\varepsilon + R_1)dT$$

$$d\Phi = (-9\varepsilon^{-2}\Phi + \varepsilon^{-2}\alpha a^2 + R_2)dT$$

$$d\Psi = (-\varepsilon^{-2}\Psi + \varepsilon^{-2}\alpha|a|^2 + \varepsilon^{-2}\alpha Z_\varepsilon^2 + R_3)dT$$

We exchange the  $a\phi$ ,  $a\psi$  and  $\varepsilon^{-1}2\alpha a Z_\varepsilon$  terms by applying Itô differentiation on these terms.

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# The rest of the proof

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The rest of the proof consists of three parts:

- ▶ Show that  $\int aZ_\varepsilon^2 dt \approx \int a\sigma dt$  (Averaging Lemma)
- ▶ Show that  $a$  is approximately  $A$
- ▶ Show that the stopping time  $\tau_*$  is long enough (i.e. bigger than a fixed time independent of  $\varepsilon$ )



# The SSH equation on unbounded domains

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The solution to the amplitude equation  $A(T)$  has values in the Sobolev space  $\mathcal{H}^\alpha$ ,  $\alpha > 1/2$  defined by

$$\mathcal{H}^\alpha := \{u \in L^2(\mathbb{R}; \mathbb{C}) : \mathcal{F}^{-1}((1 + k^2)^{\alpha/2} \mathcal{F}u) \in L^2(\mathbb{R}; \mathbb{C})\}.$$

The solution to (SSH) is approximated by

$$u(t, x) \approx \varepsilon A(\varepsilon^2 t, \varepsilon x) e^{ix} + \varepsilon \bar{A}(\varepsilon^2 t, \varepsilon x) e^{-ix} + Z_\varepsilon$$



# Problems on unbounded domains

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- ▶ No Fourier series, but Fourier transform with bands of Eigenvalues.
- ▶ SDEs for the modes feature a full linear operator instead of a scalar, which makes the exchanging of mixed products ( $a\Phi, a\Psi, \dots$ ) much more difficult.
- ▶ Bounds still depend a lot on  $A$  being in  $\mathcal{H}^{1/2+}$  which prohibits more general noise (which is at most  $\mathcal{H}^{1/2-}$ ).



Thank you.

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# Thank you for your attention!